

Heavy Flavour Hadronisation in Pythia

Peter Skands (Monash University)

- 1. Heavy-Flavour Hadronisation in the Lund Model**
- 2. Constraints**
- 3. From ee to pp**
- 4. New Theory Models in Pythia**
- 5. Some Suggestions for New Measurements**
- 6. Multiply heavy hadrons?**



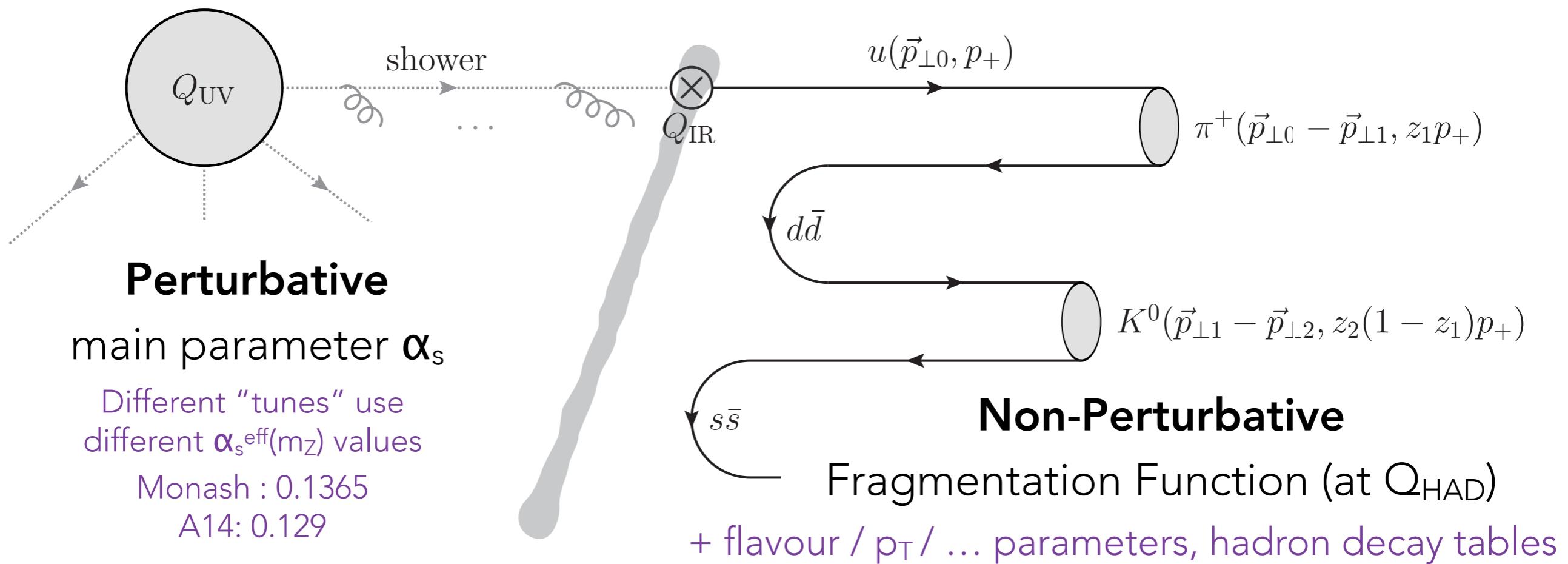
*Heavy-Flavour Hadronization in pp
& HI Collisions, CERN, March 2020*

Reminder: Fragmentation Models

Hard process (e.g., dijets) \rightarrow hard factorisation scale $Q_{\text{UV}} \sim p_{\text{Tjet}}$

Parton Showers: perturbative bremsstrahlung down to $Q_{\text{IR}} \sim 1 \text{ GeV}$

Hadronisation: confinement (+ hadron decays) at $Q_{\text{HAD}} \sim Q_{\text{IR}}$



Spectrum = combination of α_s choice & non-perturbative parameters

Flavour Composition in the Lund Model

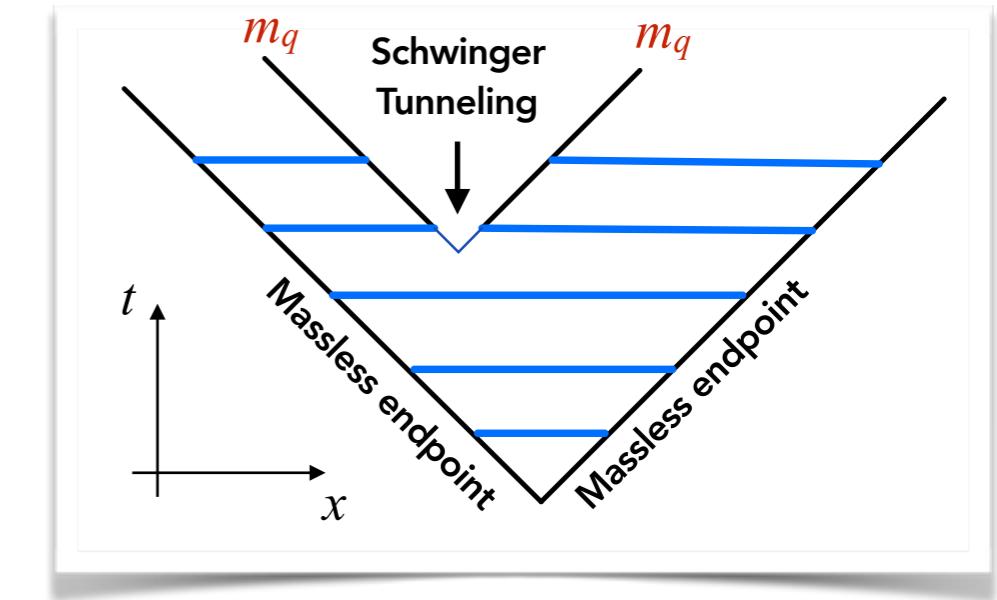
Starting point: **isolated** string in 1+1 dimensions

Tension $\kappa \sim 1 \text{ GeV/fm} \sim 0.2 \text{ GeV}^2$

String breaks by Schwinger mechanism

→ **Suppression** of strange quarks (and diquarks) $\exp\left(-\frac{m_q^2 + p_\perp^2}{\kappa}\right)$

→ StringFlav:probStoUD = 0.217



+ Spin-splitting in hadron multiplets $V/P \neq 3$

ρ/π StringFlav:mesonUDvector = 0.50

K^*/K StringFlav:mesonSvector = 0.55

D^*/D StringFlav:mesonCvector = 0.88

B^*/B StringFlav:mesonBvector = 2.2

Note: model parameters are for **primary** hadrons \neq measured ratios (feed-down)

Flavour Composition in the Lund Model

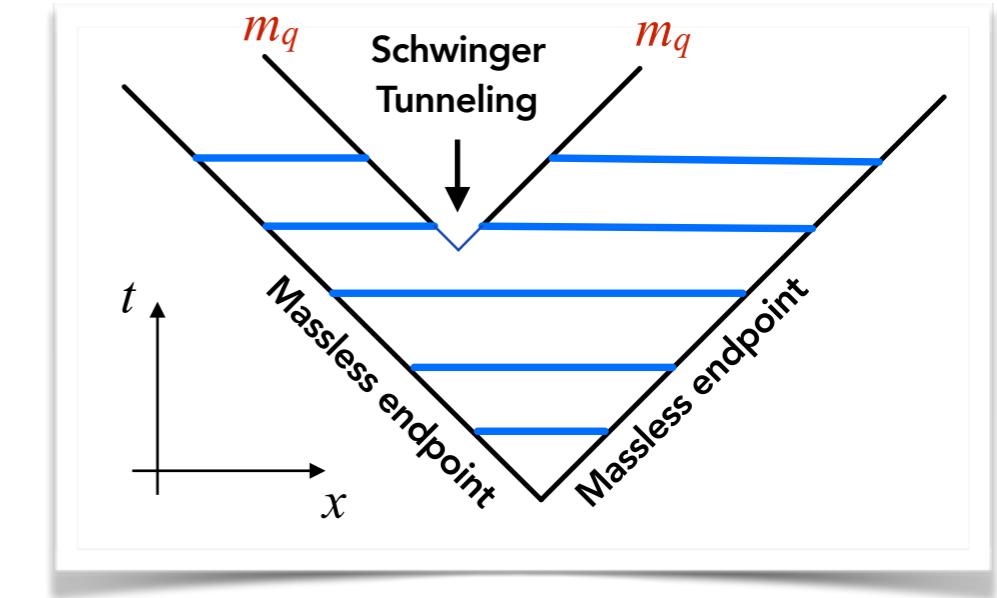
Starting point: **isolated** string in 1+1 dimensions

Tension $\kappa \sim 1 \text{ GeV/fm} \sim 0.2 \text{ GeV}^2$

String breaks by Schwinger mechanism

→ **Suppression** of strange quarks (and diquarks) $\exp\left(-\frac{m_q^2 + p_\perp^2}{\kappa}\right)$

→ StringFlav:probStoUD = 0.217



+ Spin-splitting in hadron multiplets $V/P \neq 3$

ρ/π StringFlav:mesonUDvector = 0.50

K^*/K StringFlav:mesonSvector = 0.55

D*/D StringFlav:mesonCvector = 0.88

B*/B StringFlav:mesonBvector = 2.2

Note: model parameters are for **primary** hadrons ≠ measured ratios (feed-down)

Rookie Mistake: for D*/D in the Monash tune

arXiv:1404.5630

I took the D and D* rates from separate sources ➤ wrong ratio
Should be higher ~ 1.25 - 1.5 to agree with measured values

Thanks to D. Bardhan
for pointing to this

Heavy-Flavour Endpoint Quarks

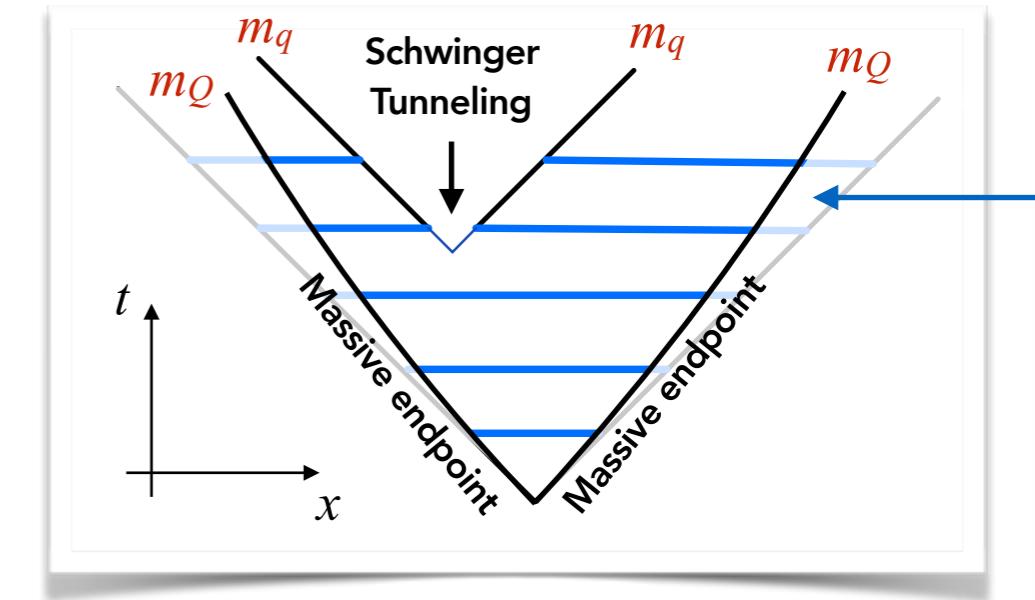
Same starting point as for massless endpoints

Tension $\kappa \sim 1 \text{ GeV/fm} \sim 0.2 \text{ GeV}^2$

String breaks by Schwinger mechanism

→ **Suppression** of strange quarks (and diquarks) $\exp\left(-\frac{m_q^2 + p_\perp^2}{\kappa}\right)$

→ StringFlav:probStoUD = 0.217



- ! Same parameters govern D_s/D , B_s/B , Λ_c/D , $\Lambda_b/B \rightarrow$ Interesting to check if D_s/D , B_s/B affected in same way in same environments where we see strangeness enhancements in light-quark sector: **multiplicity** dependence

Massive endpoints have $v < c \rightarrow$ smaller string space-time area:

→ Modified ("Lund-Bowler") FF:

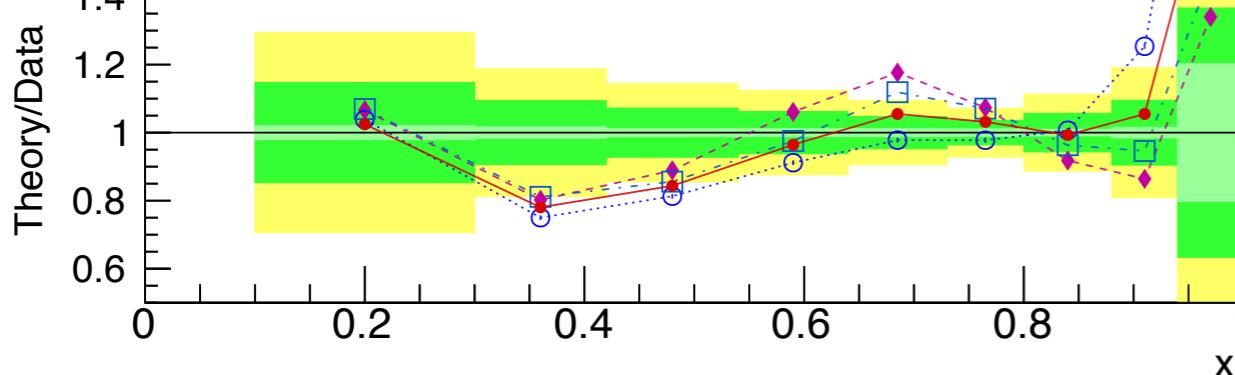
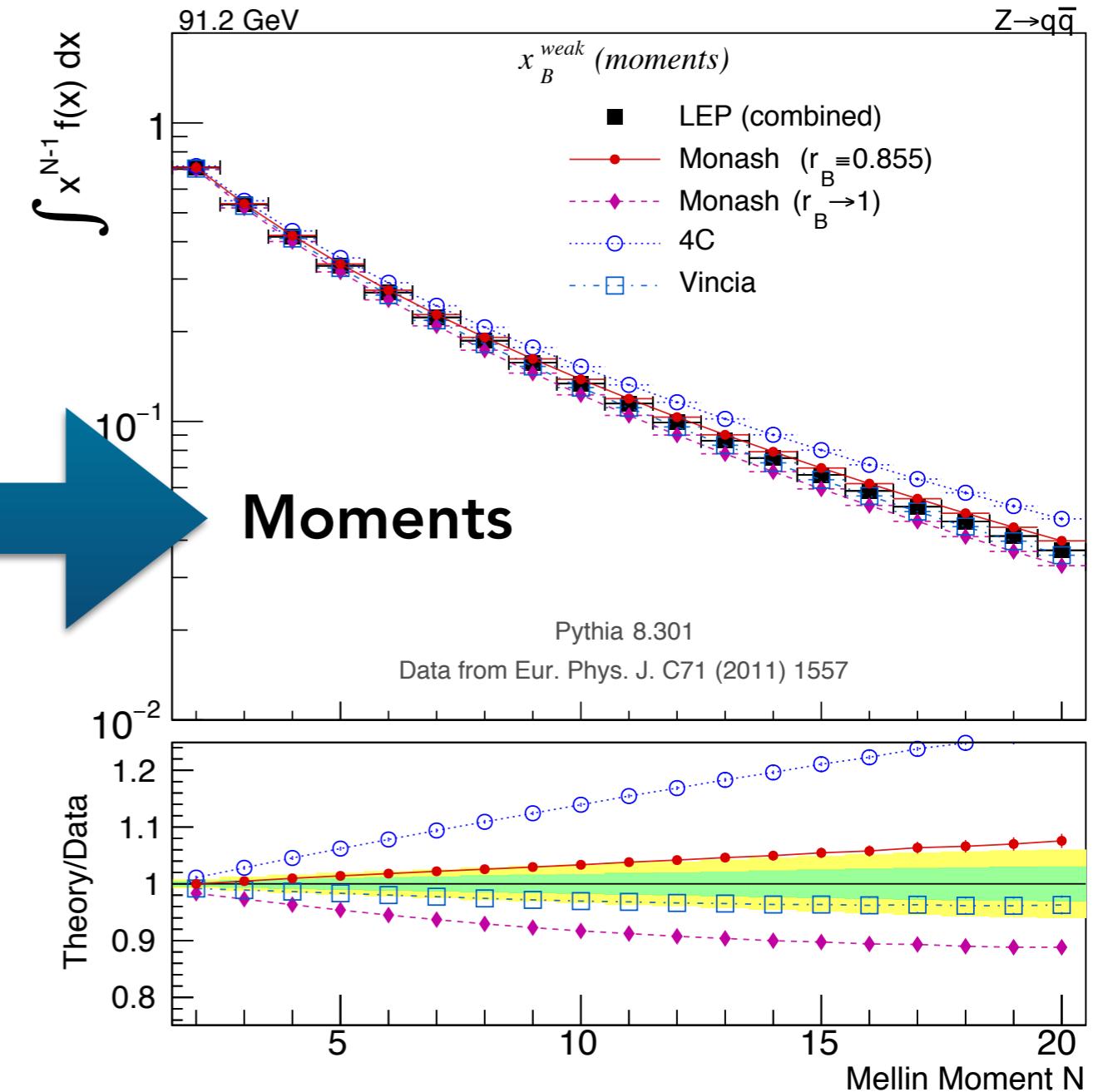
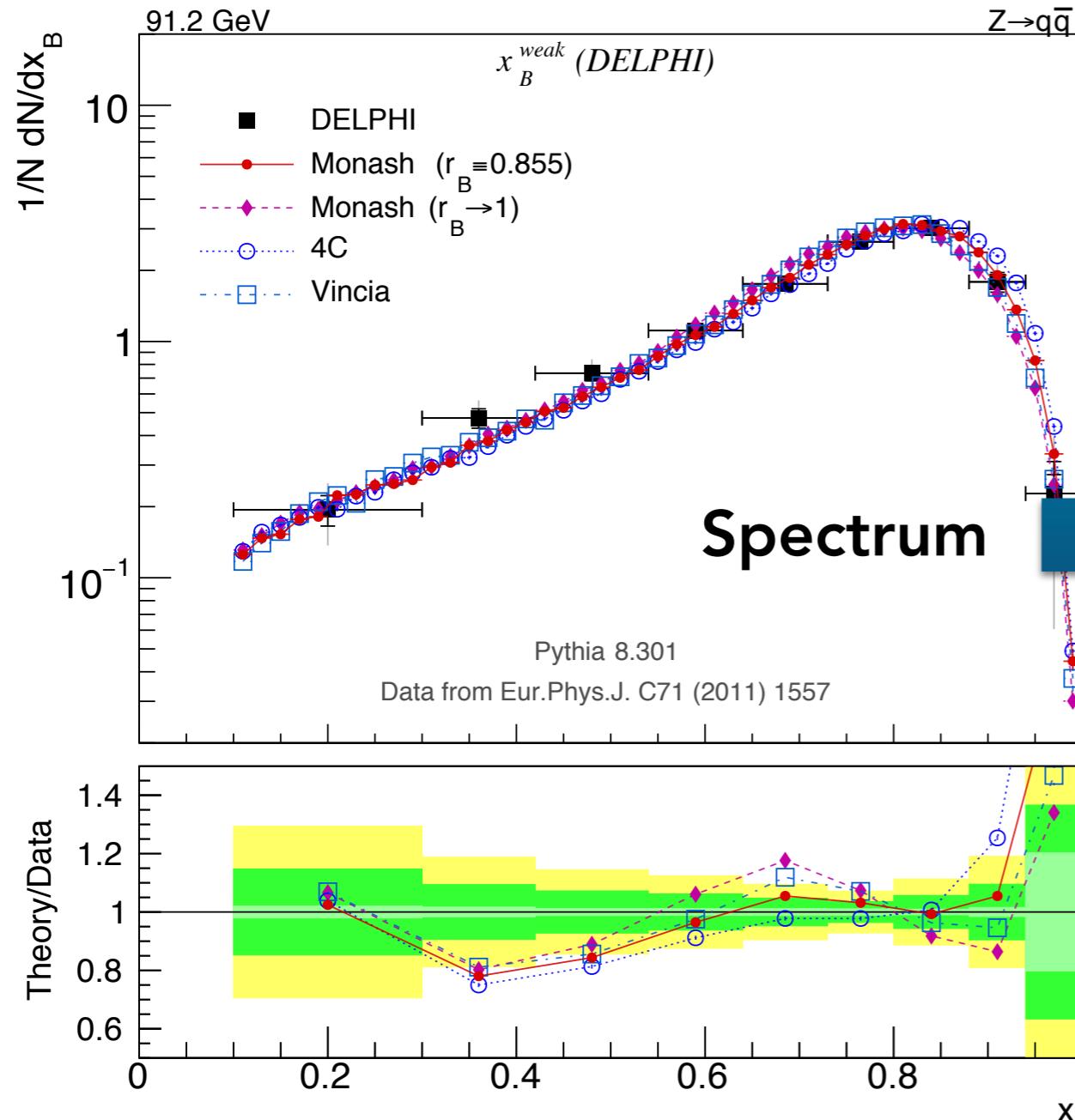
(Note: Peterson etc strictly speaking incompatible with causality in string picture)

$$\frac{(1-z)^a}{z^{1+r_Q}} \exp\left(\frac{-b m_{\perp,h}^2}{z}\right) \quad \text{with } r_b \sim r_c \sim 1$$

StringZ:rFactB = 0.855

Constraints : B Spectra

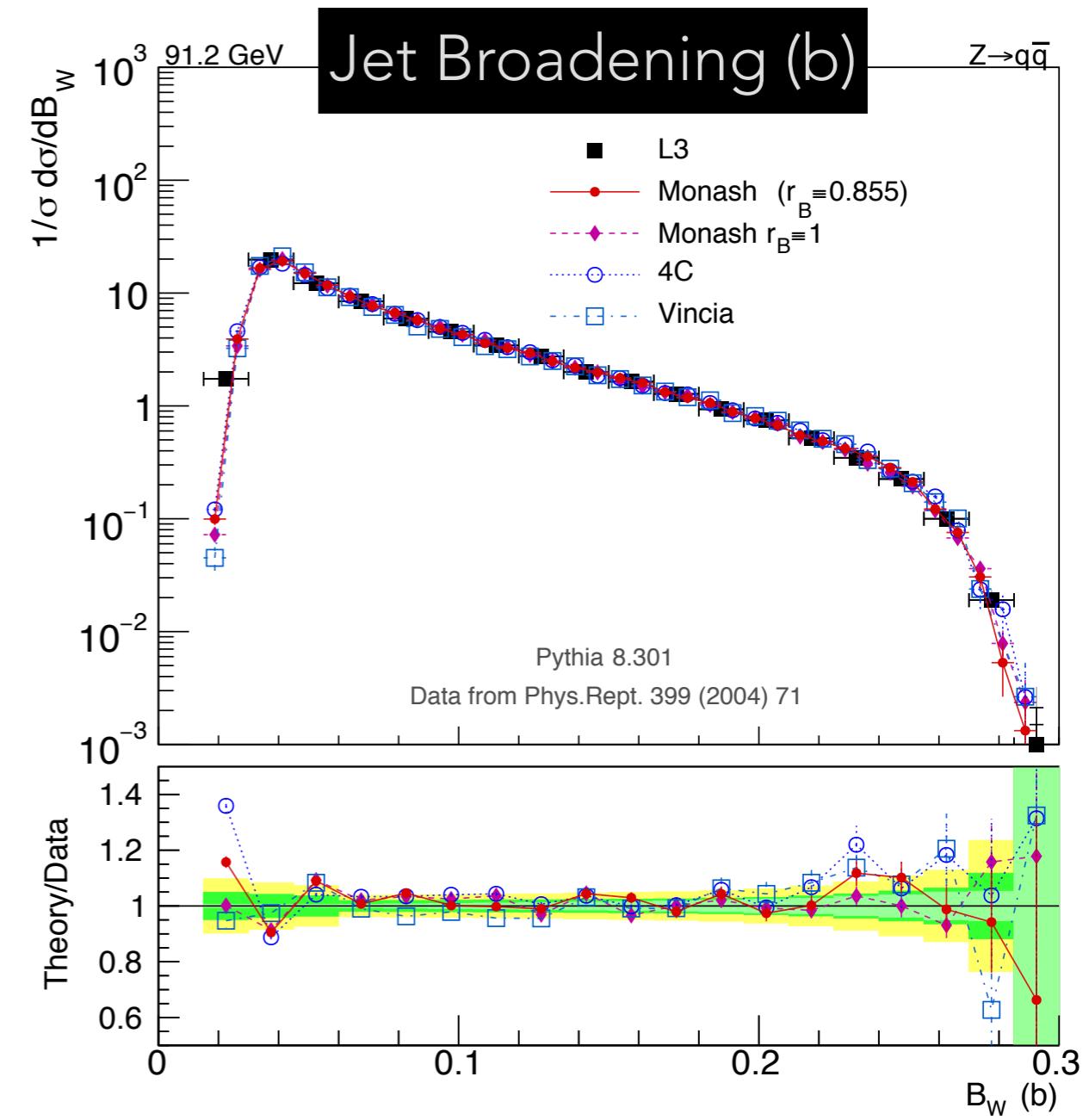
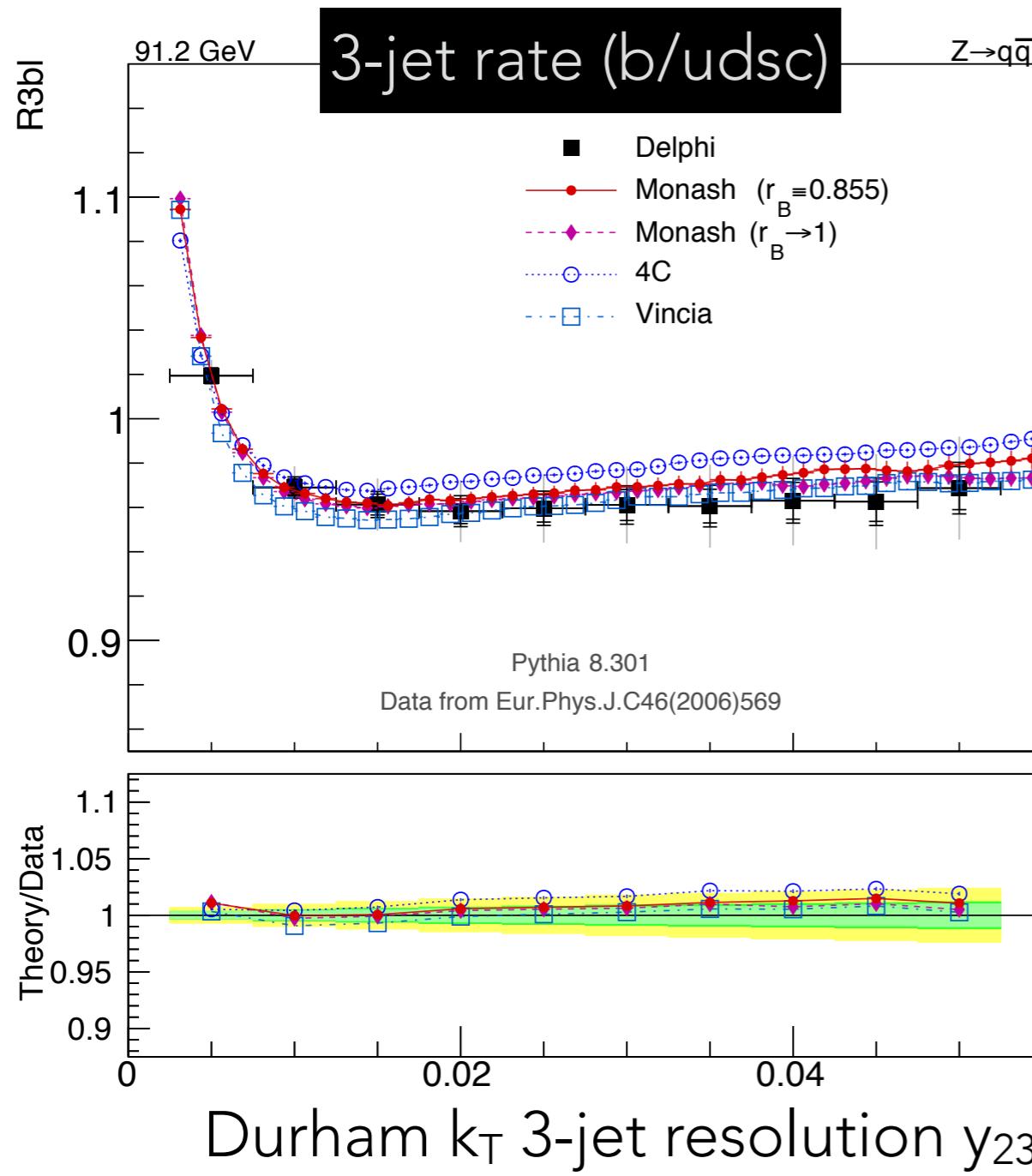
Main constraint: x_B spectra of weakly decaying B hadrons in Z decays



for details see arXiv:1404.5630 (section 2.3)

+ B-tagged Event Shapes & Jet Rates

IR safe: sensitive to α_s and b mass effects in shower + hadronisation



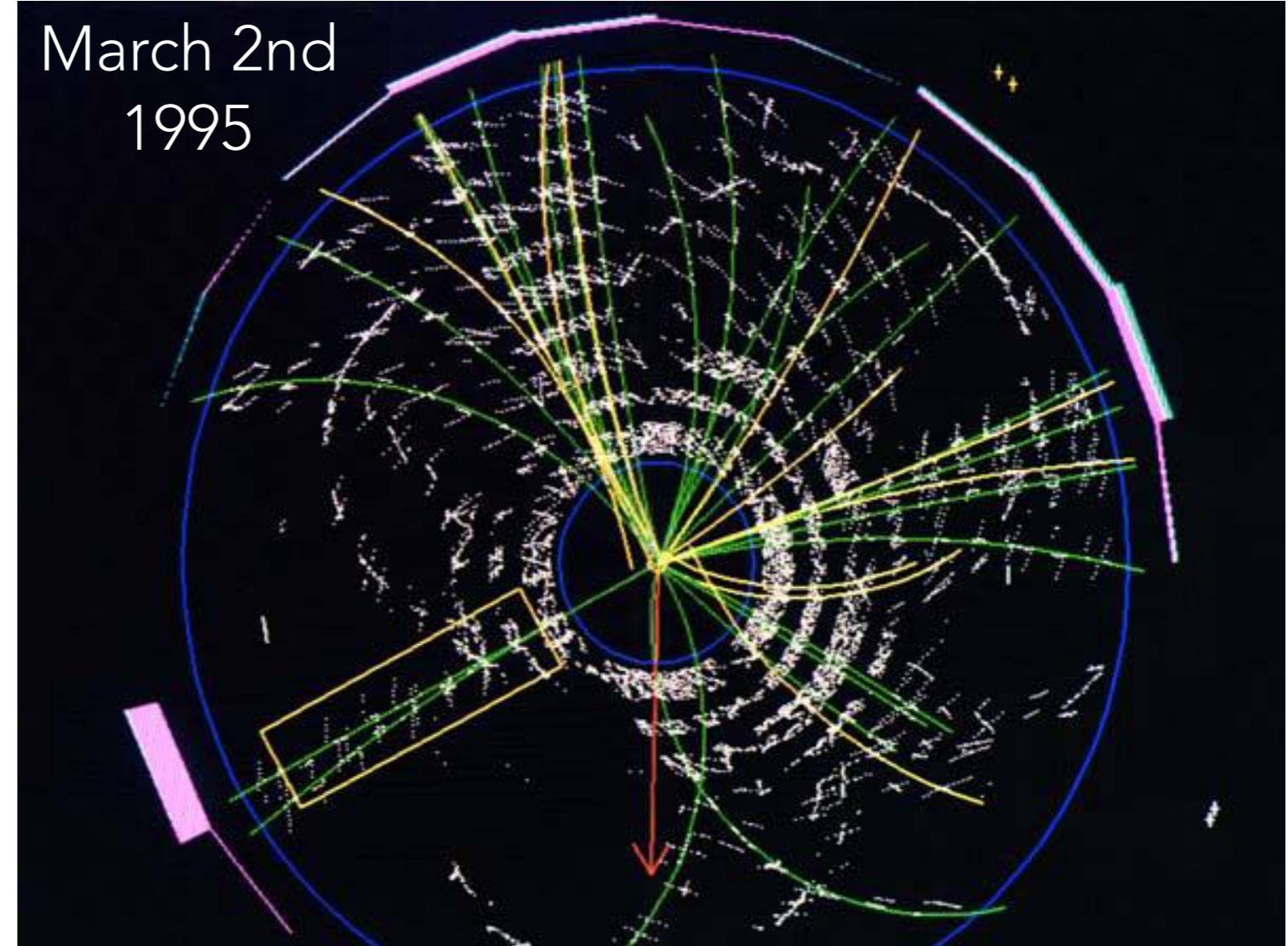
for details see arXiv:1404.5630 (section 2.3)

LHC: Top Decays > In-situ controlled B-Jet Sample?

Yesterday: 25th anniversary of the top quark discovery



$t \rightarrow bW$ provides a clean high-statistics reference sample, with a well-defined initial b-quark energy (in top CM) very similar to $Z \rightarrow bb$.



Compare $B \text{ FF}(x)$ and B hadron flavour ratios to those for inclusive b-jets, incl. any dependence on UE level (measured away from the top jets)

Note: finite top width \rightarrow “collective effects” may be suppressed in top (“early” vs “late” resonance decays)

Some Comments on b fragmentation “tuning”

Note: Monash uses “large” TimeShower:alphaSvalue = 0.1365

Regarded at least in part as making up for **NLO** K-factor for **ee→3 jets** (baseline Pythia only accurate to LO for 3 jets).

Consistent with 3-flavour $\Lambda_{\text{QCD}} \sim 0.35 \text{ GeV}$ (since we use 1-loop running)

Not guaranteed to be universal.

LHC studies tend to prefer lower effective values of α_s

E.g., A14 uses TimeShower:alphaSvalue = 0.129 (could be reinterpreted via CMW to MSbar $\alpha_s(m_Z) \sim 0.12$ so consistent with world average.)

(but I would then also change to 2-loop running to preserve Λ_{QCD} value)

E.g., a lower $\alpha_s \rightarrow$ less perturbative radiation → harder $x_b(Q_{\text{IR}})$

→ Would need to retune non-perturbative parameters (e.g., r_b) at LEP

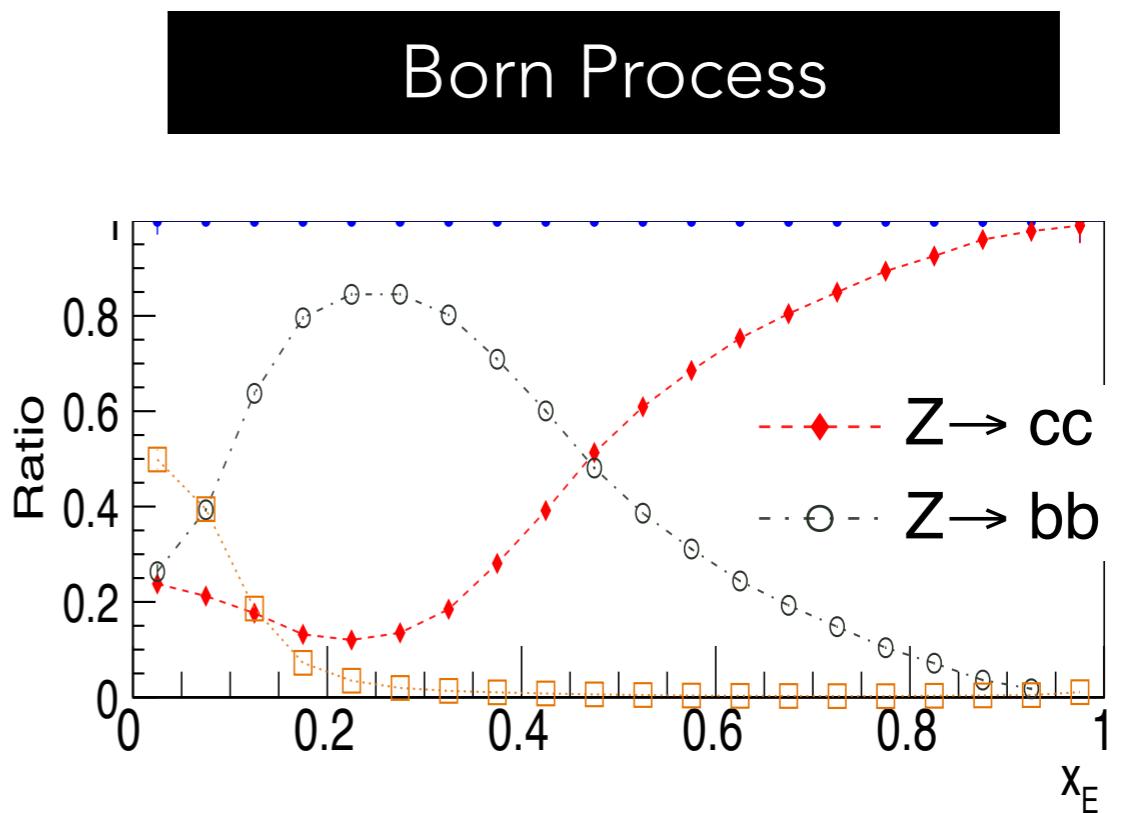
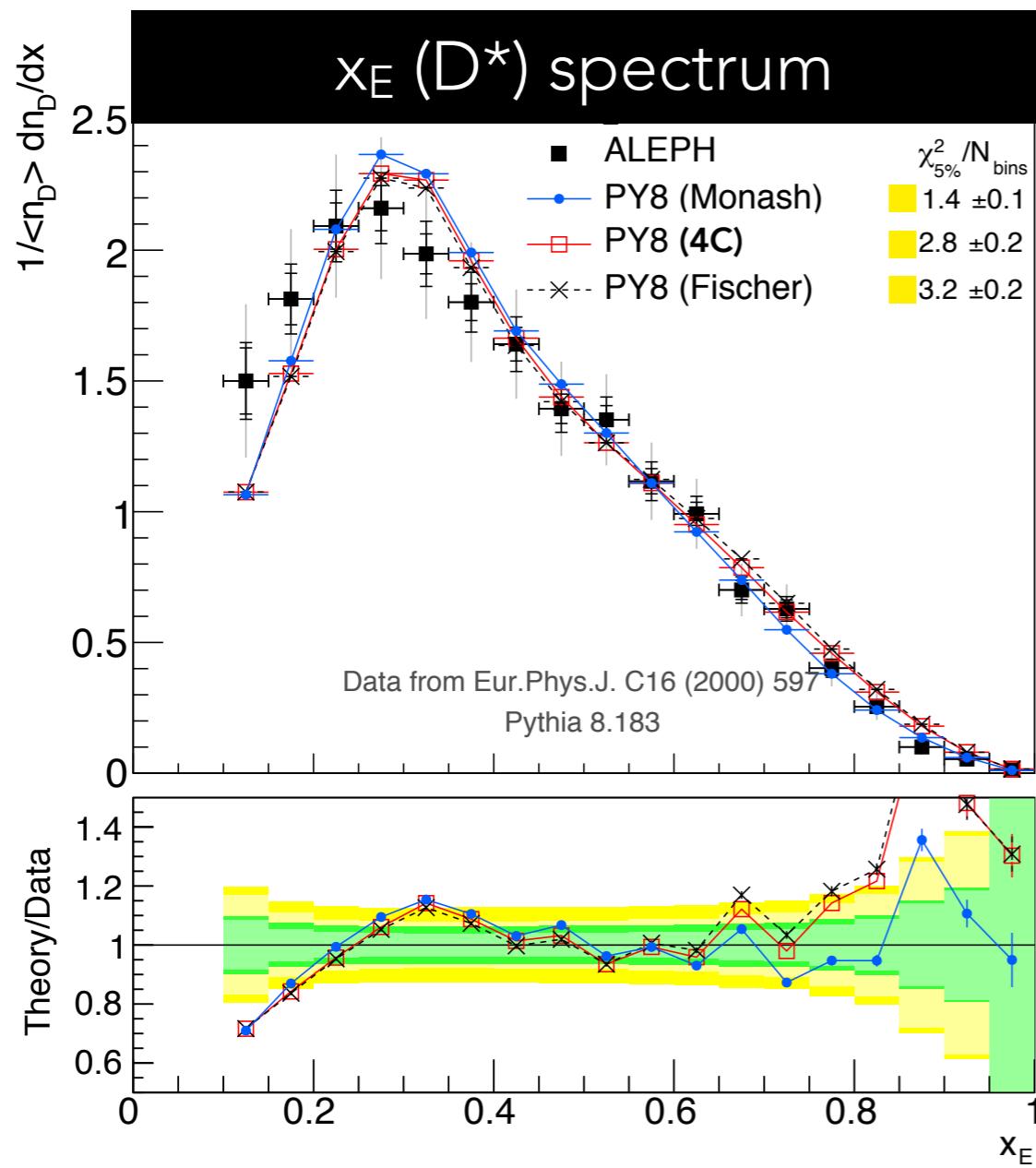
Problem: most LEP measurements are inclusive (including 3-jet events)

→ Would need 3-jet NLO merging to ensure correct 3-jet admixture.

Constraints : Charm

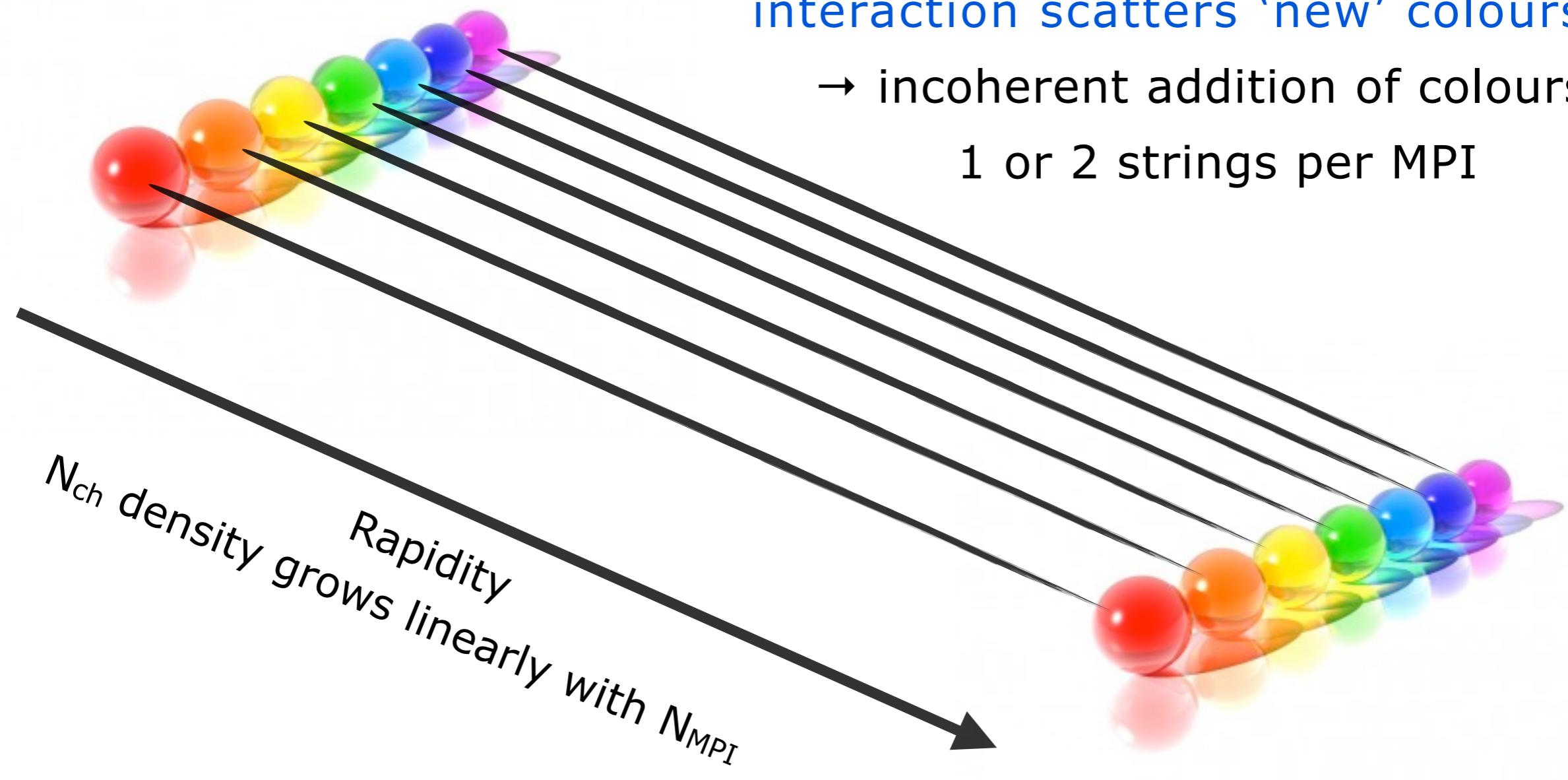
No "C-tagged" data from LEP (that I am aware of)

Monash tune only used a single D^* spectrum (ALEPH) $\rightarrow r_c$



Actual charm fragmentation not seen very clearly. For $x < 0.5$, the inclusive D^* spectrum is dominated by B decays

From ee to pp: multiple parton interactions (MPI)

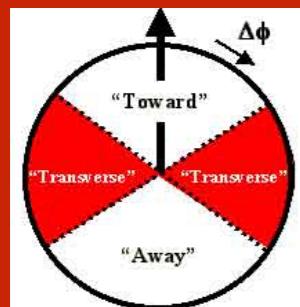


Simple, clean, factorized picture ...

WRONG!

Anticipated already in first Pythia MPI model (*Sjöstrand & van Zijl, 1987*)

“CR” parameter = probability for MPI to just generate “kinks” on hard-process colour structure, rather than new strings of their own



Tevatron $\langle N_{ch} \rangle$ and $\langle \Sigma p_T \rangle$ in “Transverse” UE region
Required ~ 100% CR (*Rick Field, “Tune A”, 2002*)

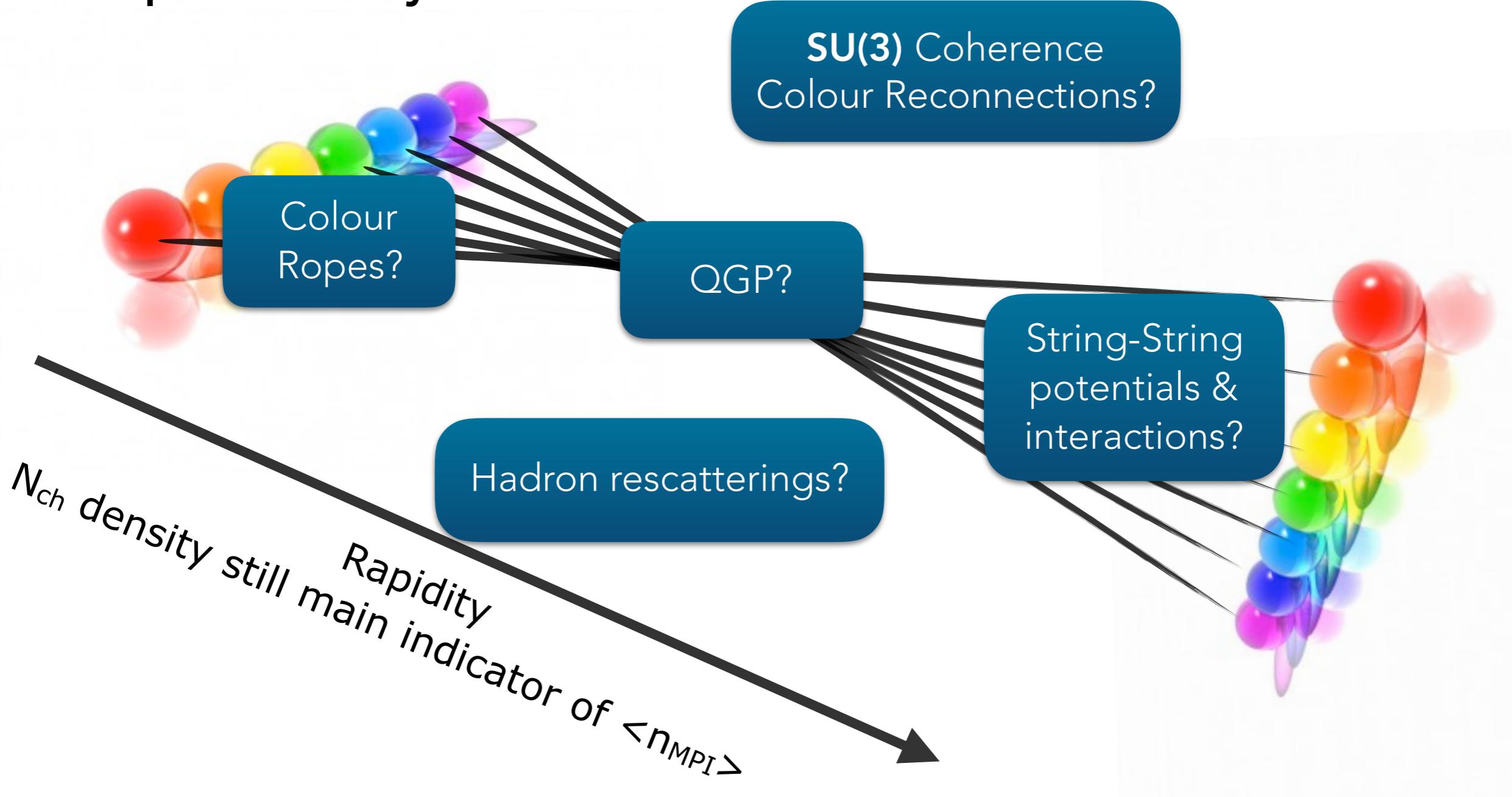


► Not a small effect, then ...

+ Many new measurements and discoveries from LHC (& RHIC)
(e.g., CMS ridge, ALICE strangeness vs N_{ch} , ...)

The MPI are all within a proton radius of each other (in pp)

The picture today



Brief Summary of “New” Theory Models in Pythia

QCD-inspired CR ColourReconnection:mode = 1 Christiansen, Skands, JHEP 1508 (2015) 003

Stochastically sample subleading- N_C connections according to **SU(3) weights** and choose among possibilities (incl colour- ϵ ones) based on **string-length minimisation**.

► **some flow effects & additional baryons** (incl multiply-heavy); no extra strangeness

Ropes & Shoving Bierlich, Gustafson, Lönnblad, Tarasov, JHEP 1503 (2015) 148 Bierlich, Gustafson, Lönnblad, PLB 779 (2018) 58

Ropes: allow QCD charges to combine into higher representations: 6, 10, 15, 21, 28, ... with **higher string tensions** (Casimir scaling) ► **more strangeness & more baryons**

Shoving: explicit dynamical model of repulsion between different strings/ropes ► **flow**

Thermodynamical String fragmentation Fischer, Sjöstrand, JHEP 1701 (2017) 140

+ Much ongoing work ...

Hadronic Rescattering (Sjöstrand+Utheim)

HI extensions (Angantyr, PISTA) & extensions with **UrQMD** (Bierlich et al.)

Interacting Strings: momentum-space alternative to ropes+shoving (Duncan+Skands)

Back to basics: fragmentation of a single string: early / out-of-equilibrium, and thermal effects. **Time-varying string tension** out soon. + other variants? E.g., **UCLA model?**

Some Suggestions for New Measurements

Want to disentangle $\langle p_T \rangle$, $\langle \text{strangeness} \rangle$, $\langle \text{baryons} \rangle$, $\langle N_{\text{ch}} \rangle$
 $\langle \zeta \rangle$ ($\langle \varsigma \rangle$) $\langle \mathcal{B} \rangle$

E.g., CR and “flow” increase $\langle p_T \rangle$ without (directly) affecting $\langle \zeta \rangle$

“Baryonic” CR can increase $\langle \mathcal{B} \rangle$

Higher tensions/temperatures: correlated $\langle p_T \rangle$, $\langle \zeta \rangle$, and $\langle \mathcal{B} \rangle$

Some Simple Questions:

How **local** are the $\langle \zeta \rangle$ and $\langle \mathcal{B} \rangle$ enhancement mechanisms?

How far in phase space is **nearest anti-strange / anti-baryon?**

For different values of N_{ch} density, $p_{T\text{B}}$ or $p_T(\text{b-jet})$, and ζ density

E.g., heavy-flavour tag, say $B_s \rightarrow$ know the endpoint flavour \rightarrow look for nearest anti-strange quark.

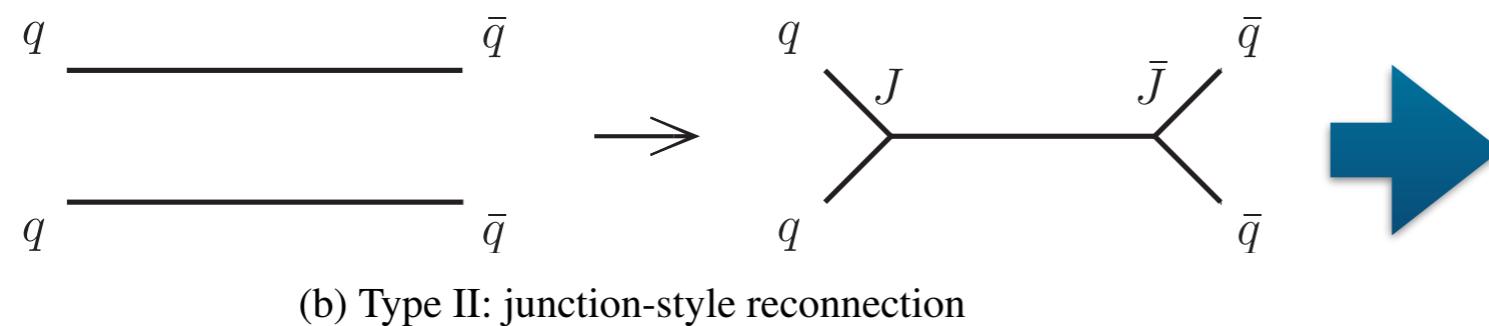
What is the distance in p_T ? in rapidity (along z / along b-jet)? in ΔR ?

How do the HF fractions depend on **event multiplicity?**

Heavy-Flavour Baryons

Example: QCD-inspired CR

Allows “junction reconnections”, e.g.:



For the parameters used in that study,

Λ_c/D^+ increased by factor 2

Λ_b/B^+ by factor 3

+ potentially larger changes for $\Sigma_{c,b}^{(*)}$

Christiansen, Skands, JHEP 1508 (2015) 003

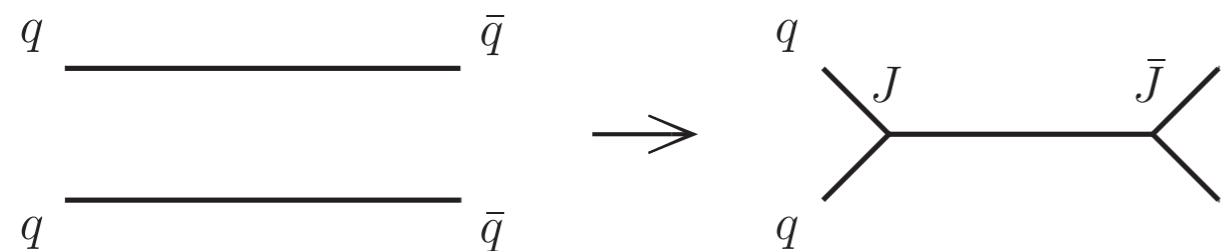
ColourReconnection:mode = 1 = 0

Particle	$N_{\text{par}}/N_{\text{events}}$	
D^+	$5.3 \cdot 10^{-2}$	$6.5 \cdot 10^{-2}$
Λ_c^+	$1.2 \cdot 10^{-2}$	$6.6 \cdot 10^{-3}$
Σ_c^{++}	$1.3 \cdot 10^{-2}$	$5.4 \cdot 10^{-4}$
Σ_c^+	$1.5 \cdot 10^{-2}$	$5.2 \cdot 10^{-4}$
Σ_c^0	$1.3 \cdot 10^{-2}$	$5.1 \cdot 10^{-4}$
Σ_c^{*++}	$2.2 \cdot 10^{-3}$	$9.5 \cdot 10^{-4}$
Σ_c^{*+}	$2.4 \cdot 10^{-3}$	$9.4 \cdot 10^{-4}$
Σ_c^{*0}	$2.2 \cdot 10^{-3}$	$9.1 \cdot 10^{-4}$
ccq^7	$2.1 \cdot 10^{-4}$	$1.0 \cdot 10^{-7}$
B^+	$1.6 \cdot 10^{-3}$	$2.3 \cdot 10^{-3}$
Λ_b^0	$8.2 \cdot 10^{-4}$	$3.9 \cdot 10^{-4}$
Σ_b^+	$9.5 \cdot 10^{-4}$	$3.1 \cdot 10^{-5}$
Σ_b^0	$1.0 \cdot 10^{-3}$	$3.7 \cdot 10^{-5}$
Σ_b^-	$9.4 \cdot 10^{-4}$	$3.2 \cdot 10^{-5}$
Σ_b^{*+}	$9.5 \cdot 10^{-4}$	$3.1 \cdot 10^{-5}$
Σ_b^{*0}	$1.0 \cdot 10^{-3}$	$3.7 \cdot 10^{-5}$
Σ_b^{*-}	$9.4 \cdot 10^{-4}$	$3.2 \cdot 10^{-5}$
bcq^7	$1.8 \cdot 10^{-5}$	0
bbq^7	$1.1 \cdot 10^{-6}$	0

Heavy-Flavour Baryons

Example: QCD-inspired CR

Allows “junction reconnections”, e.g.:

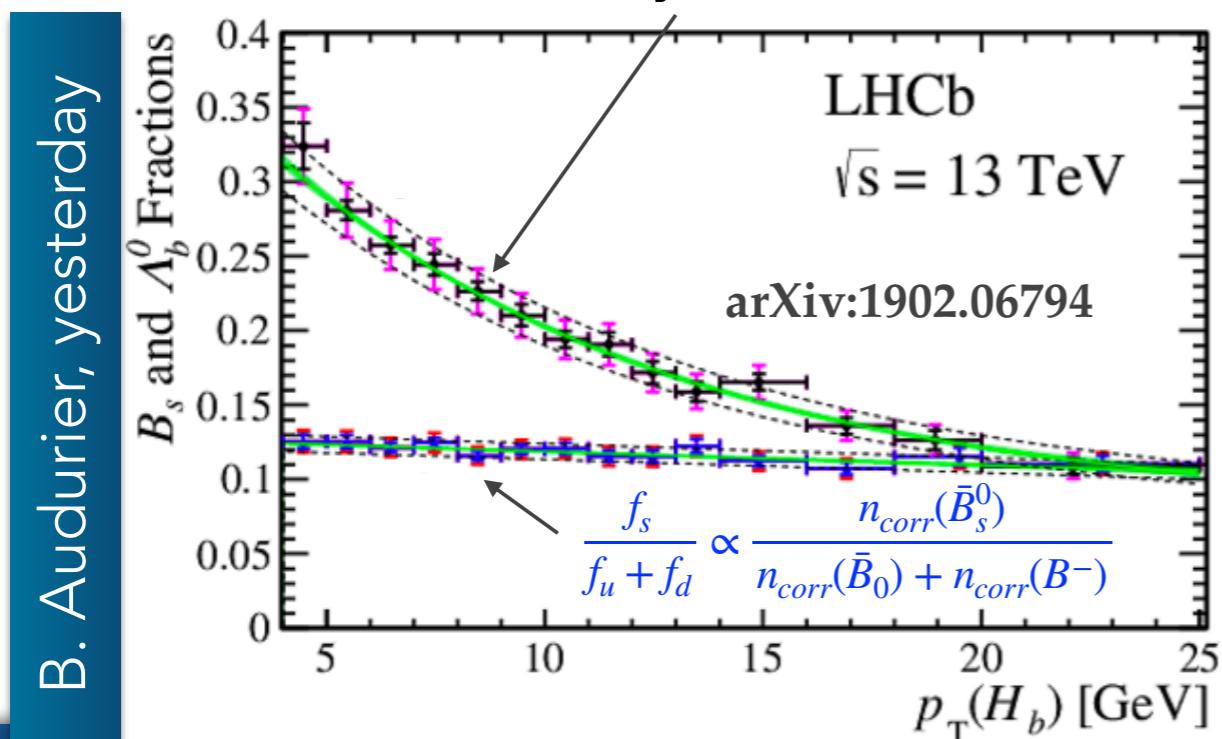


(b) Type II: junction-style reconnection

For the parameters used in that study,

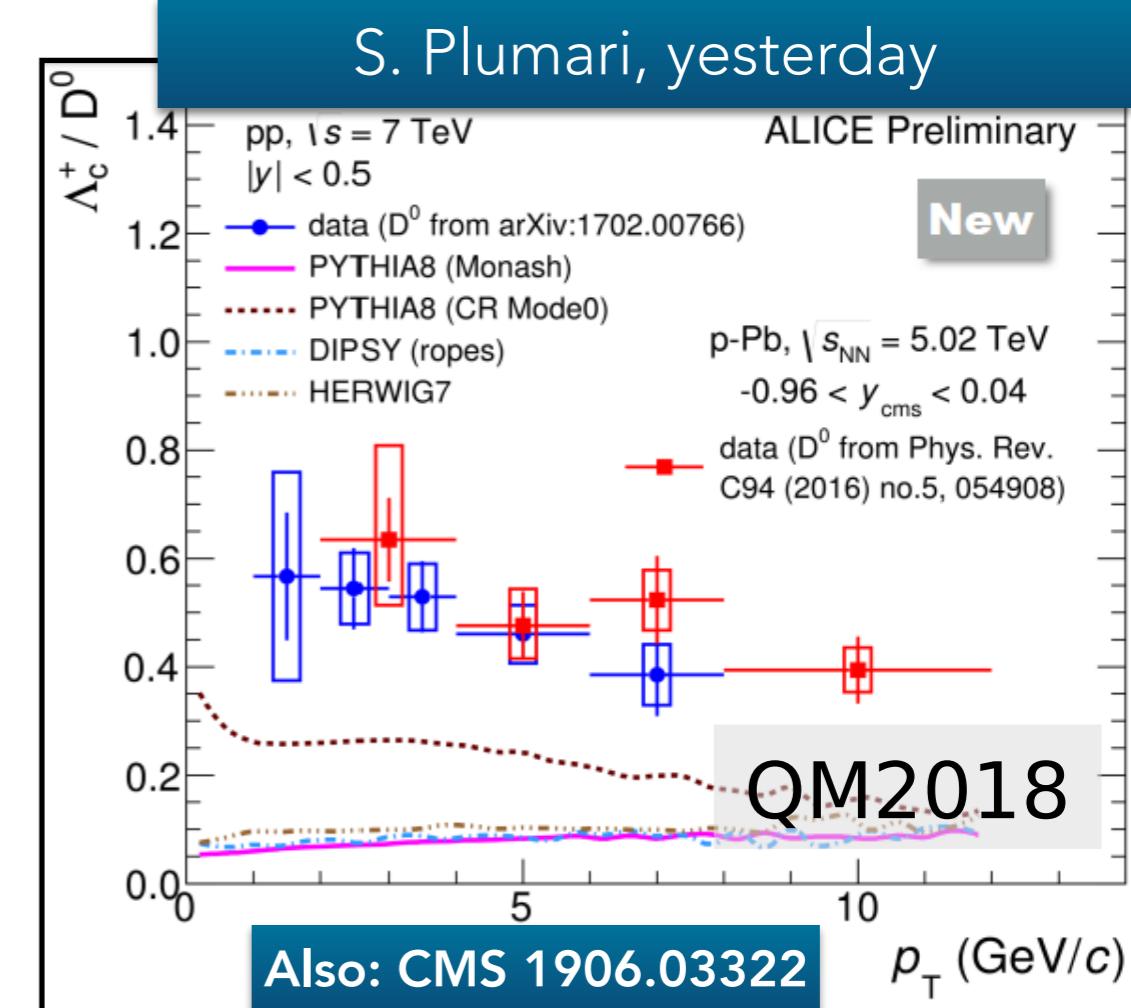
Λ_c/D^+ increased by factor 2

Λ_b/B^+ by factor 3



Christiansen, Skands, JHEP 1508 (2015) 003
ColourReconnection:mode = 1 = 0

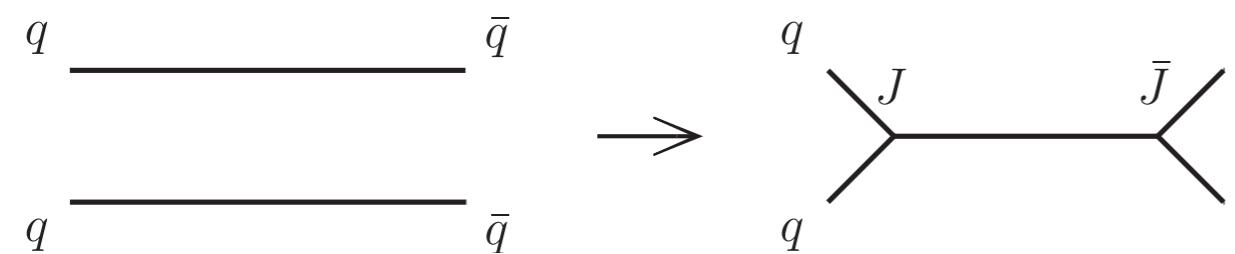
Particle	$N_{\text{par}}/N_{\text{events}}$
D^+	$5.3 \cdot 10^{-2}$
Λ_c^+	$1.2 \cdot 10^{-2}$
Σ_c^{++}	$1.3 \cdot 10^{-2}$
Σ_c^+	$1.5 \cdot 10^{-2}$
Σ^0	$1.3 \cdot 10^{-2}$



Heavy-Flavour Baryons

Example: QCD-inspired CR

Allows “junction reconnections”, e.g.:



(b) Type II: junction-style reconnection

For the parameters used in that study,

Λ_c/D^+ increased by factor 2

Λ_b/B^+ by factor 3

+ potentially larger changes for $\Sigma_{c,b}^{(*)}$

Generically expect dependence on multiplicity \rightarrow Measure $\langle B/M \rangle(N_{ch})$?

(Should be true for ropes, hydro, ... too)

+ baryon-antibaryon rapidity dependence?

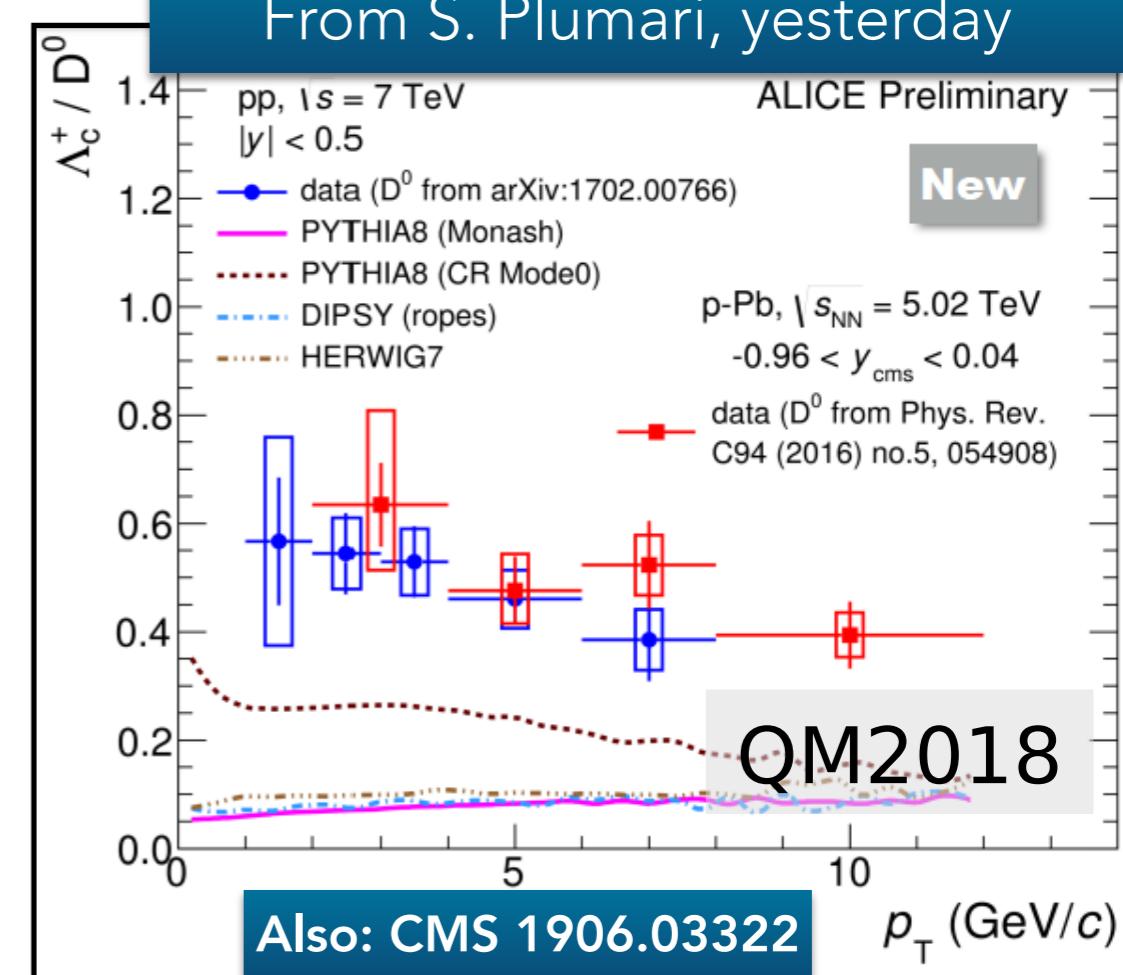
First step, e.g., ALICE D meson associated track multiplicities in arXiv:1910.14403

Christiansen, Skands, JHEP 1508 (2015) 003

ColourReconnection:mode = 1 = 0

Particle	$N_{\text{par}}/N_{\text{events}}$
D^+	$5.3 \cdot 10^{-2}$
Λ_c^+	$1.2 \cdot 10^{-2}$
Σ_c^{++}	$1.3 \cdot 10^{-2}$
Σ_c^+	$1.5 \cdot 10^{-2}$
Σ^0	$1.3 \cdot 10^{-2}$

From S. Plumari, yesterday



(Some) LHCb measurements

B_s/B^+ vs event kinematics

From $B_s \rightarrow J/\psi \varphi$ & $B^+ \rightarrow J/\psi K^+$

No dependence on p_{LB} , η_B

Decreasing trend with p_{TB}



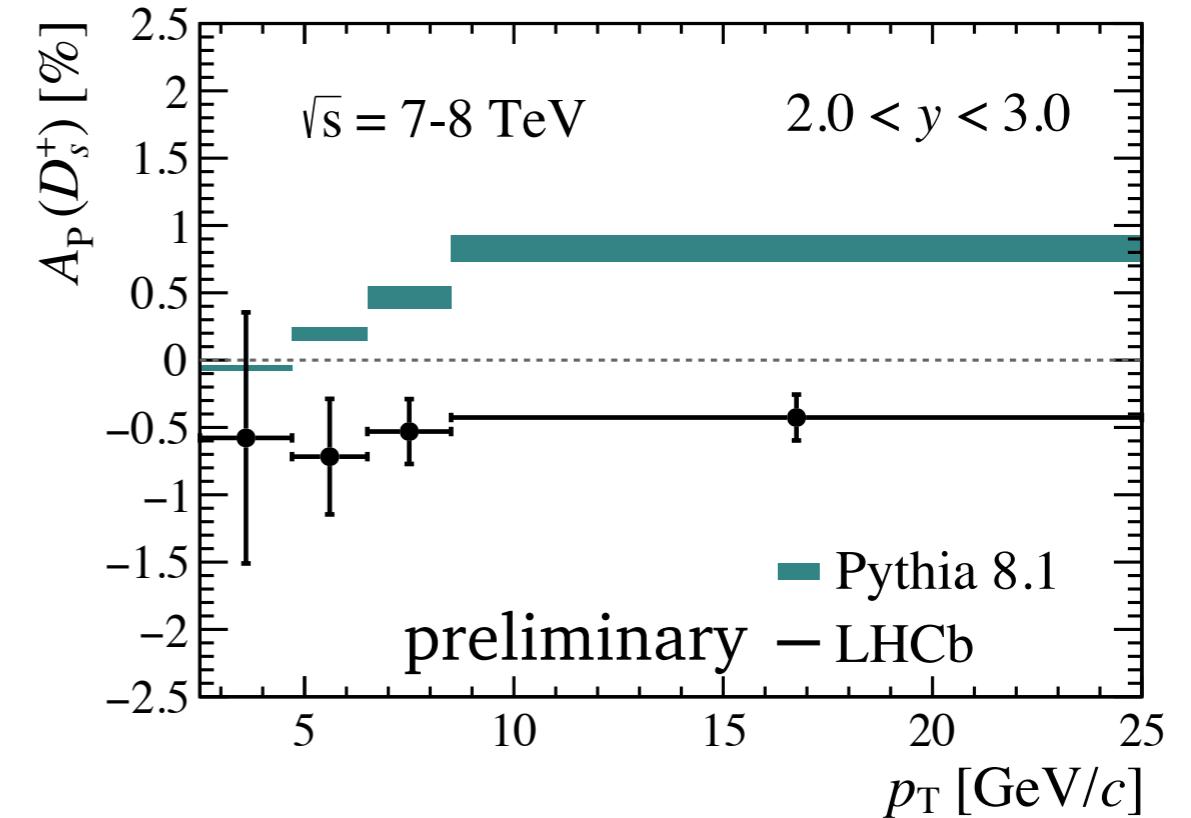
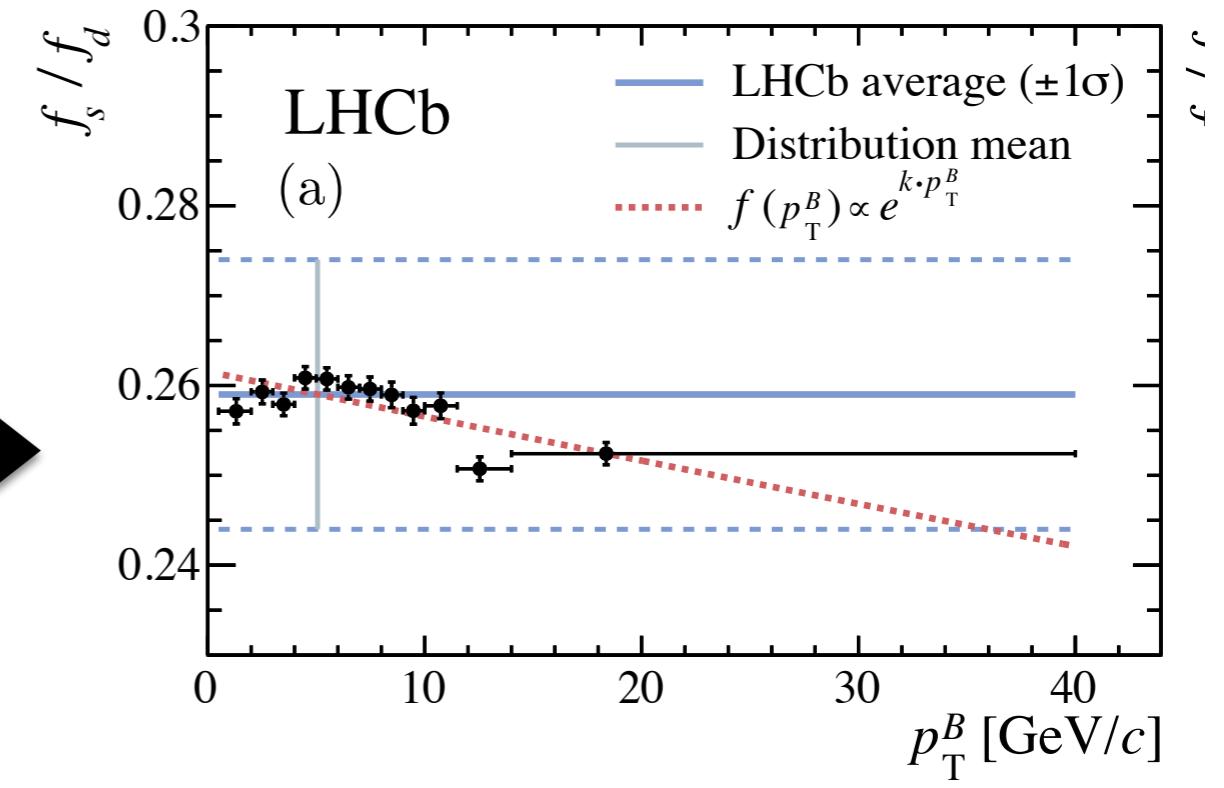
Would be highly interesting to see vs event multiplicity / associated track multiplicity

D_s asymmetry (S. Klaver, Moriond 2018)

$$\frac{\sigma(pp \rightarrow D_s^+) - \sigma(pp \rightarrow D_s^-)}{\sigma(pp \rightarrow D_s^+) + \sigma(pp \rightarrow D_s^-)}$$

Strong p_T dependence in Pythia, not seen in data

High p_T ➤ Coherence effect?



Multiply Heavy Hadrons?

Heavy flavours produced perturbatively, not in string/cluster breakups;

So why would multiply heavy hadrons be interesting as soft probes?

Because they also **probe the confinement field** in unique ways (colour- ϵ_{ijk})

E.g., the Ξ_{cc} has been measured

[LHCb-PAPER-2019-037](#)

Does its rate vary with associated track density?

Christiansen, Skands, JHEP 1508 (2015) 003

Particle	ColourReconnection:mode = 1			= 0
	New CR model ($N_{\text{par}}/N_{\text{events}}$)			Old CR model
	string	junction	all	$N_{\text{par}}/N_{\text{events}}$ (all)
D^+	$5.3 \cdot 10^{-2}$	0	$5.3 \cdot 10^{-2}$	$6.5 \cdot 10^{-2}$
Λ_c^+	$4.0 \cdot 10^{-3}$	$7.9 \cdot 10^{-3}$	$1.2 \cdot 10^{-2}$	$6.6 \cdot 10^{-3}$
Σ_c^{++}	$2.7 \cdot 10^{-4}$	$1.3 \cdot 10^{-2}$	$1.3 \cdot 10^{-2}$	$5.4 \cdot 10^{-4}$
Σ_c^+	$2.5 \cdot 10^{-4}$	$1.5 \cdot 10^{-2}$	$1.5 \cdot 10^{-2}$	$5.2 \cdot 10^{-4}$
Σ_c^0	$2.5 \cdot 10^{-4}$	$1.3 \cdot 10^{-2}$	$1.3 \cdot 10^{-2}$	$5.1 \cdot 10^{-4}$
Σ_c^{*++}	$5.1 \cdot 10^{-4}$	$1.7 \cdot 10^{-3}$	$2.2 \cdot 10^{-3}$	$9.5 \cdot 10^{-4}$
Σ_c^{*+}	$4.9 \cdot 10^{-4}$	$1.9 \cdot 10^{-3}$	$2.4 \cdot 10^{-3}$	$9.4 \cdot 10^{-4}$
Σ_c^{*0}	$4.8 \cdot 10^{-4}$	$1.7 \cdot 10^{-3}$	$2.2 \cdot 10^{-3}$	$9.1 \cdot 10^{-4}$
ccq^7	0	$2.1 \cdot 10^{-4}$	$2.1 \cdot 10^{-4}$	$1.0 \cdot 10^{-7}$
B^+	$1.6 \cdot 10^{-3}$	0	$1.6 \cdot 10^{-3}$	$2.3 \cdot 10^{-3}$
Λ_b^0	$1.9 \cdot 10^{-4}$	$6.3 \cdot 10^{-4}$	$8.2 \cdot 10^{-4}$	$3.9 \cdot 10^{-4}$
Σ_b^+	$1.1 \cdot 10^{-5}$	$9.3 \cdot 10^{-4}$	$9.5 \cdot 10^{-4}$	$3.1 \cdot 10^{-5}$
Σ_b^0	$1.2 \cdot 10^{-5}$	$1.0 \cdot 10^{-3}$	$1.0 \cdot 10^{-3}$	$3.7 \cdot 10^{-5}$
Σ_b^-	$1.1 \cdot 10^{-5}$	$9.3 \cdot 10^{-4}$	$9.4 \cdot 10^{-4}$	$3.2 \cdot 10^{-5}$
Σ_b^{*+}	$1.1 \cdot 10^{-5}$	$9.3 \cdot 10^{-4}$	$9.5 \cdot 10^{-4}$	$3.1 \cdot 10^{-5}$
Σ_b^{*0}	$1.2 \cdot 10^{-5}$	$1.0 \cdot 10^{-3}$	$1.0 \cdot 10^{-3}$	$3.7 \cdot 10^{-5}$
Σ_b^{*-}	$1.1 \cdot 10^{-5}$	$9.3 \cdot 10^{-4}$	$9.4 \cdot 10^{-4}$	$3.2 \cdot 10^{-5}$
bcq^7	0	$1.8 \cdot 10^{-5}$	$1.8 \cdot 10^{-5}$	0
bbq^7	0	$1.1 \cdot 10^{-6}$	$1.1 \cdot 10^{-6}$	0

Note: the baryon "predictions" depend on poorly constrained model parameters; highlight measurement sensitivity

To discuss: **observables to tell apart ...**

CR: longitudinal (1D) strings + transverse boosts: flow-like effects

No $\langle \zeta \rangle$ enhancement; low velocity dispersions relative to common boosts

Additional tracers: multiply heavy baryons (will at least ➤ constraints!)

Ropes etc: longitudinal (1D) strings with higher effective tensions

Strangeness enhancement + higher $\langle p_T \rangle$, but still "1D"

➤ rank ordering, const dN/dy ?

Shoving etc.: Longitudinal strings with transverse repulsions

1D "rank" still relevant for $\langle \mathcal{B} \rangle$, $\langle \zeta \rangle$, and (local) p_T conservation ➤ correlations?

+ higher tensions? Is $\langle p_T \rangle$ correlated or anti-correlated with $\langle \zeta \rangle$, $\langle \mathcal{B} \rangle$?

Thermal/Statistical systems: 3D systems with higher effective T

Very high dispersions, 3D.

Quantum number and p_T conservation not ordered in "rank" at all?

Extra Slides

What a strange world we live in, said Alice [to the queen of hearts]

We wanted to know if “violent” collision events produced higher-strength fields.

Smoking gun would be a higher fraction of strange particles being produced

(higher-strength fields \Rightarrow more energy per “space-time volume” \Rightarrow easier to produce higher-mass quark-antiquark pairs)

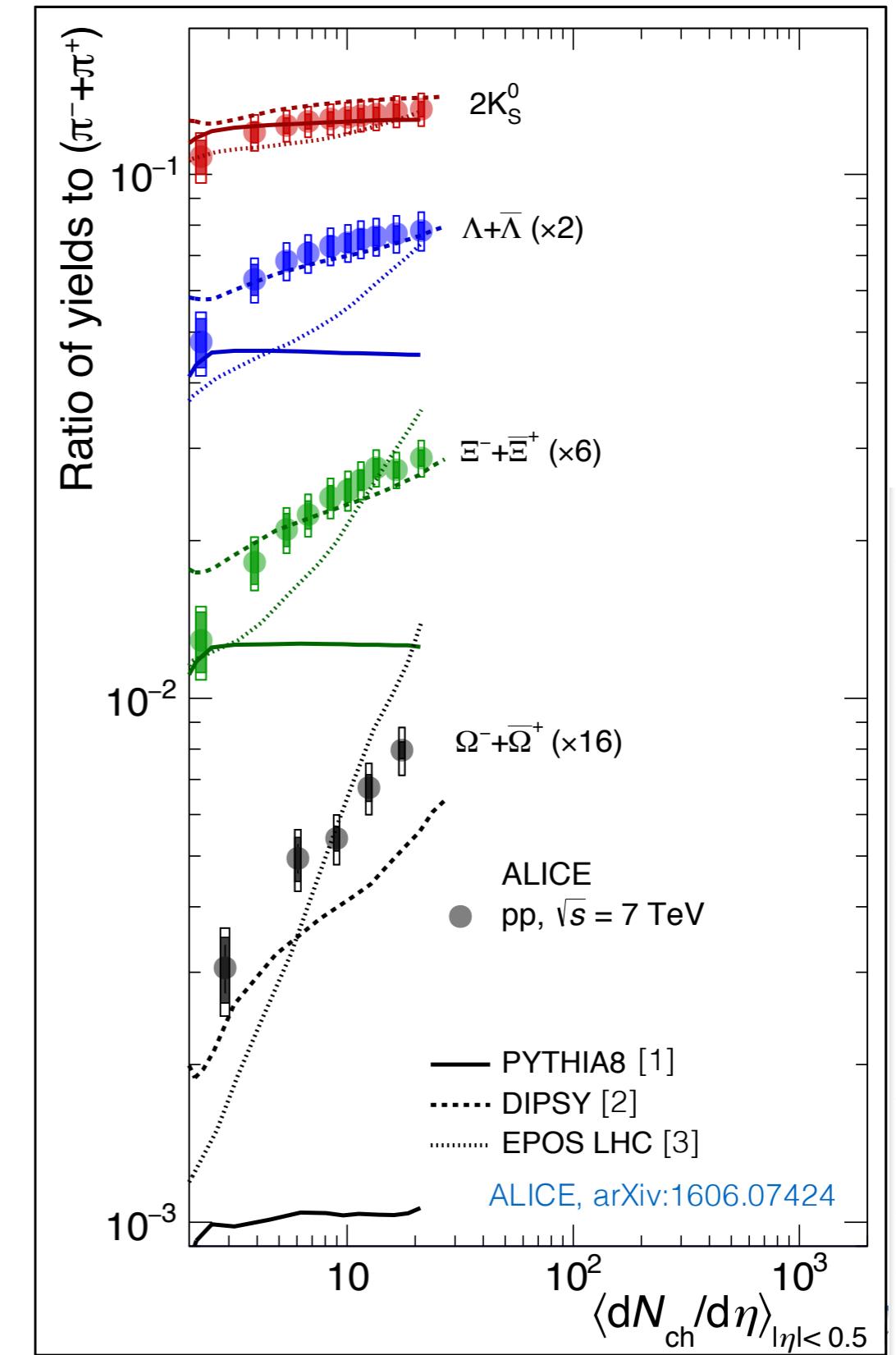
Jackpot!

Now working on models in which nearby fragmenting fields interact with each other.

Interactions between QCD strings!

Higher tensions + repulsion effects \blacktriangleright modifications in high-density environments

(Competing idea: the whole thing turns into a near-perfect liquid which gets heated up.)



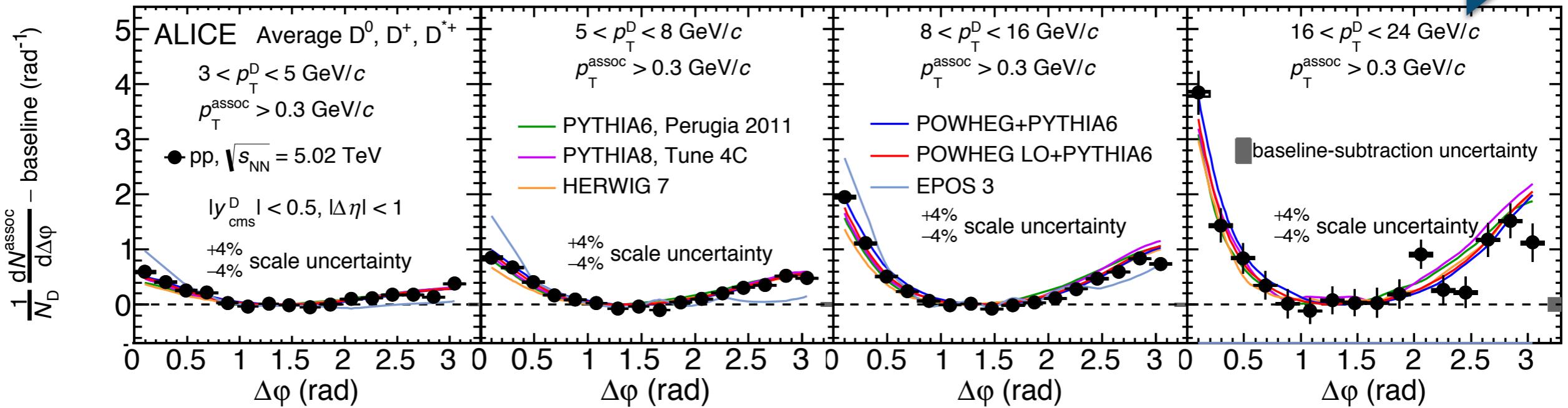
D meson associated tracks (ALICE)

arXiv:1910.14403v1

Soft

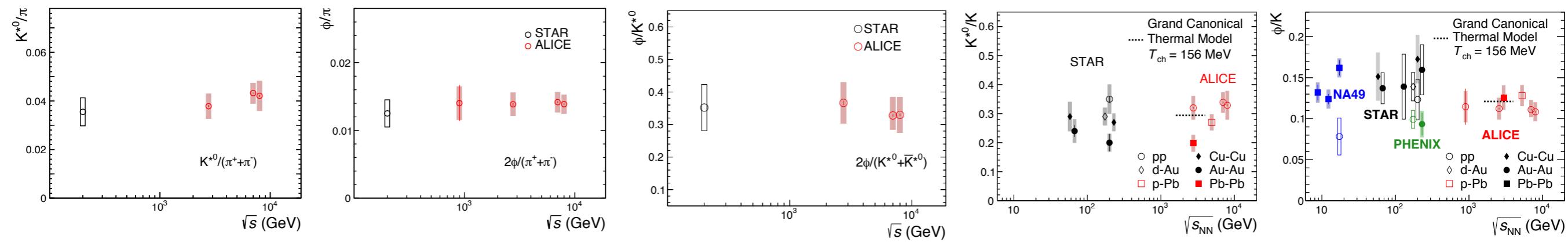
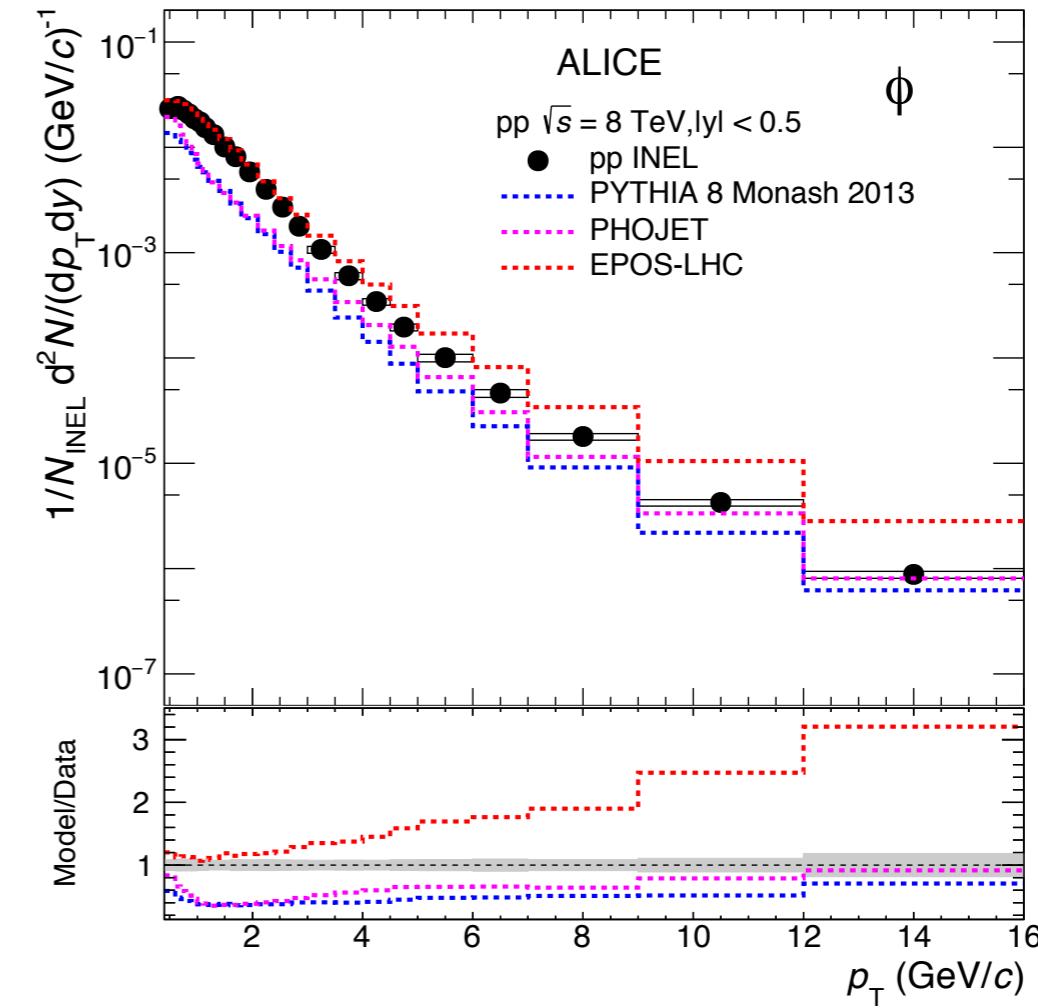
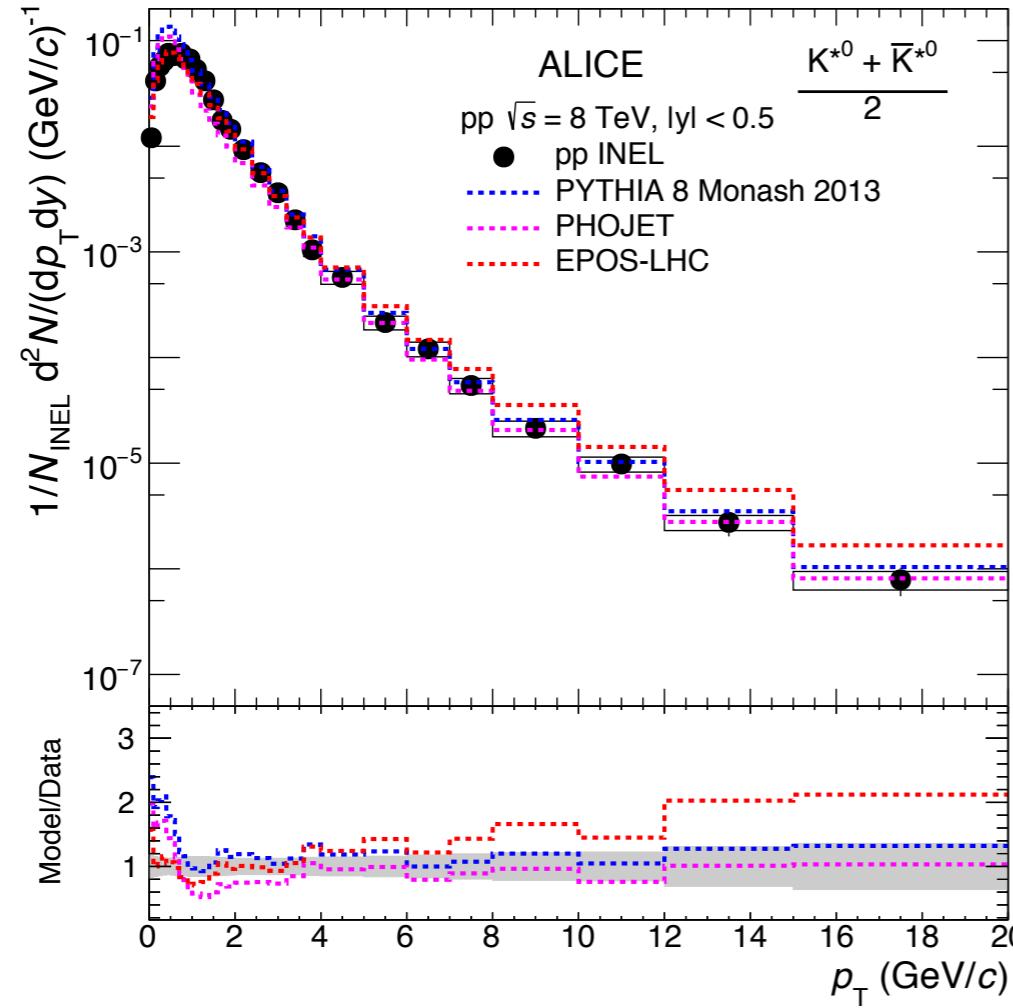
D meson pT

Hard



K* and ϕ (ALICE)

arXiv:1910.14410v1



NB: n_{ch} dependence measured separately, in arXiv:1910.14397

K* and ϕ multiplicity dependence (ALICE)

arXiv:1910.14397v1

