Photoproduction and Ultraperipheral collisions in Pythia 8

Oxford theory seminar

Ilkka Helenius

May 2nd, 2024









Outline

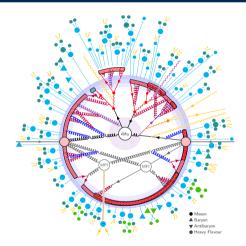
Pythia 8: A general purpose event generator

- Latest release 8.311 (March 2024)
- A new physics manual for 8.3

[SciPost Phys. Codebases 8-r8.3 (2022)]

Outline

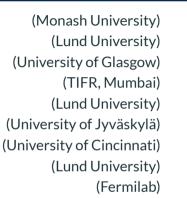
- 1. Pythia 8 basics
- 2. Structure of real photons
- 3. Photoproduction in e+p
- 4. Ultraperipheral collisions (UPCs)
- 5. Summary & Outlook



[figure by P. Skands]

Pythia Collaboration

- Javira Altmann
- Christian Bierlich
- Naomi Cooke
- Nishita Desai
- Leif Gellersen
- Ilkka Helenius
- Philip Ilten
- Leif Lönnblad
- Stephen Mrenna
- Christian Preuss (University of Wuppertal)
- Torbjörn Sjöstrand (Lund University)
- Peter Skands (Monash University/Oxford)
- Marius Utheim (University of Jyväskylä)
- Rob Verheyen (University College London)



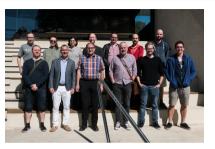


[Pythia meeting in Monash 2019]

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(Monash University) (Lund University) (University of Glasgow) (TIFR, Mumbai) (Lund University) (University of Jyväskylä) (University of Cincinnati) (Lund University) (Fermilab)



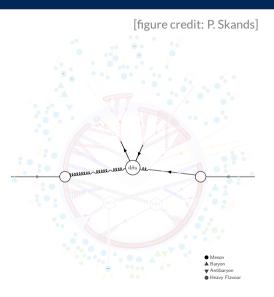
[Pythia meeting in Monash 2019]

- Spokesperson
- Codemaster
- Webmaster

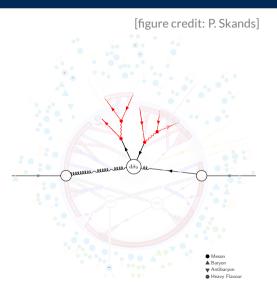
https://pythia.org authors@pythia.org

Classify event generation in terms of "hardness"

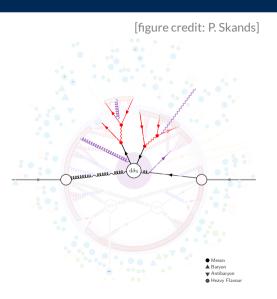
1. Hard Process (here $t\bar{t}$)



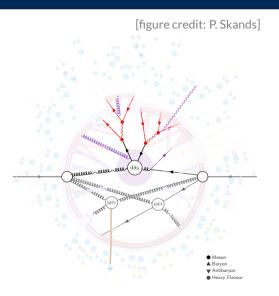
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- 2. Resonance decays (t, Z, \ldots)



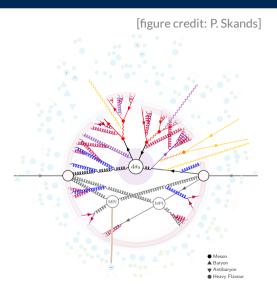
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- 3. Matching, Merging and matrix-element corrections



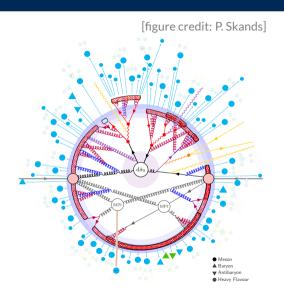
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- 4. Multiparton interactions



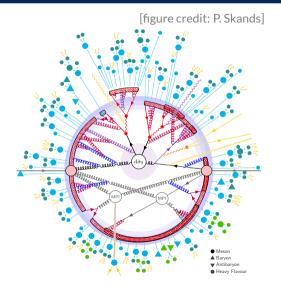
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- 7. Decays, Rescattering



Structure of real photons

Electron-proton collisions

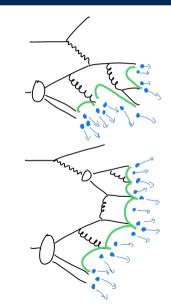
Classified in terms photon virtuality Q^2

Deep inelastic scattering (DIS)

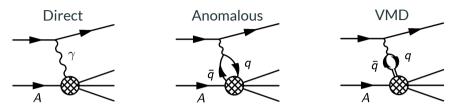
- High virtuality, $Q^2 > {\sf a}\, few\, GeV^2$
- Lepton scatters off a parton by exchanging a highly virtual photon

Photoproduction (PhP)

- Low virtuality, $Q^2 \rightarrow 0 \, GeV^2$
- Factorize photon flux, evolve γp system
- Photon may fluctuate into a hadronic state, resolved in the interaction
- Hard scale μ provided by the final state



Photon structure at $Q^2 \approx 0 \text{ GeV}^2$



Partonic structure of resolved (anom. + VMD) photon encoded in photon PDFs

$$f_i^{\gamma}(\mathbf{x}_{\gamma},\mu^2) = f_i^{\gamma,\text{dir}}(\mathbf{x}_{\gamma},\mu^2) + f_i^{\gamma,\text{anom}}(\mathbf{x}_{\gamma},\mu^2) + f_i^{\gamma,\text{VMD}}(\mathbf{x}_{\gamma},\mu^2)$$

- $f_i^{\gamma, \operatorname{dir}}(\mathbf{x}_{\gamma}, \mu^2) = \delta_{i\gamma}\delta(1 \mathbf{x}_{\gamma})$
- $f_i^{\gamma,\text{anom}}(x_{\gamma},\mu^2)$: Perturbatively calculable
- $f_i^{\gamma, \text{VMD}}(x_{\gamma}, \mu^2)$: Non-perturbative, fitted or vector-meson dominance (VMD)

Factorized cross section

$$\mathrm{d}\sigma^{\gamma \mathsf{A} \to \mathsf{k} \mathsf{l} + \mathsf{X}} = f_i^{\gamma}(\mathsf{x}_{\gamma}, \mu^2) \otimes f_j^{\mathsf{A}}(\mathsf{x}_{\mathsf{p}}, \mu^2) \otimes \mathrm{d}\sigma^{ij \to \mathsf{k} \mathsf{l}}$$

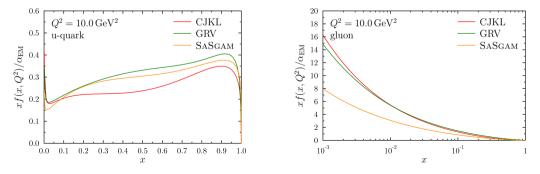
PDFs for resolved photons

DGLAP equation for photons

• Additional term due to $\gamma \rightarrow q\overline{q}$ splittings

$$\frac{\partial f_i^{\gamma}(\mathbf{x}, Q^2)}{\partial \log(Q^2)} = \frac{\alpha_{\text{em}}}{2\pi} \mathbf{e}_i^2 \mathbf{P}_{i\gamma}(\mathbf{x}) + \frac{\alpha_{\text{s}}(Q^2)}{2\pi} \sum_i \int_x^1 \frac{\mathrm{d}z}{z} \, \mathbf{P}_{ij}(z) \, f_j(x/z, Q^2)$$

where $P_{i\gamma}(x) = 3(x^2 + (1 - x)^2)$ for quarks, 0 for gluons (LO)



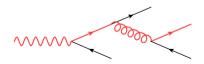
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ISR probability based on DGLAP evolution

• Add a term corresponding to $\gamma \rightarrow q\overline{q}$ to (conditional) ISR probability

$$\mathrm{d}\mathcal{P}_{a\leftarrow b} = \frac{\mathrm{d}Q^2}{Q^2} \frac{\alpha_{\rm s}}{2\pi} \frac{x' f_a^{\gamma}(x',Q^2)}{x f_b^{\gamma}(x,Q^2)} P_{a\rightarrow bc}(z) \,\mathrm{d}z + \frac{\mathrm{d}Q^2}{Q^2} \frac{\alpha_{\rm em}}{2\pi} \frac{e_b^2 \, P_{\gamma\rightarrow bc}(x)}{f_b^{\gamma}(x,Q^2)}$$

- Corresponds to ending up to the beam photon during evolution
 - \Rightarrow Parton originated from the point-like (anomalous) part of the PDFs
 - No further ISR or MPIs below the scale of the splitting
 - Implemented for the default Simple Shower in Pythia 8



Multiparton interactions (MPIs) with resolved photons

- MPIs from 2 \rightarrow 2 QCD cross sections

$d\mathcal{P}_{MPI}$ _	1	$\mathrm{d}\sigma^{2 ightarrow 2}$
$dp_T^2 =$	$\sigma_{\rm nd}(\sqrt{s})$	dp_T^2

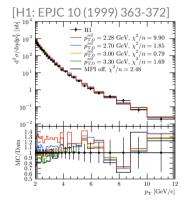
 $\sigma_{\rm nd}(\sqrt{\rm s})$ is the non-diffractive cross section

• Partonic cross section diverges at $p_{\rm T} \rightarrow 0$

 \Rightarrow Introduce a screening parameter p_{T0}

$$\frac{\mathsf{d}\sigma^{2\to2}}{\mathsf{d}p_{\mathsf{T}}^2} \propto \frac{\alpha_{\mathsf{s}}(p_{\mathsf{T}}^2)}{p_{\mathsf{T}}^4} \rightarrow \frac{\alpha_{\mathsf{s}}(p_{\mathsf{T}0}^2 + p_{\mathsf{T}}^2)}{(p_{\mathsf{T}0}^2 + p_{\mathsf{T}}^2)^2}$$

- Energy-dependent parametrization: $p_{T0}(\sqrt{s}) = p_{T0}^{ref}(\sqrt{s}/\sqrt{s_{ref}})^{\alpha}$
- Number of interactions: $\langle n \rangle = \sigma_{\rm int}(p_{\rm T0})/\sigma_{\rm nd}$



- Use H1 data to (re-)tune parameter(s)
- $\langle W_{\gamma p}
 angle pprox 200 \, {\rm GeV}$

Photoproduction in e+p

Photoproduction in electron-proton collisions

Direct processes

• Convolute photon flux f_{γ} with proton PDFs f_i^p and $d\hat{\sigma}$

 $\mathrm{d}\sigma^{\mathrm{ep}\to kl+X} = f^{\mathrm{e}}_{\gamma}(\mathbf{x}, \mathsf{Q}^2) \, \otimes \, f^{\mathrm{p}}_{j}(\mathbf{x}_{\mathrm{p}}, \mu^2) \, \otimes \, \mathrm{d}\hat{\sigma}^{\gamma j \to kl}$

• Generate FSR and ISR for proton side

Resolved processes

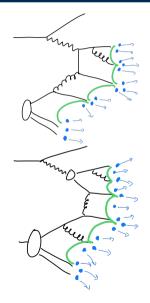
• Convolute also with photon PDFs

 $\mathrm{d}\sigma^{\mathrm{ep}\to kl+X} = f_{\gamma}^{\mathrm{e}}(\mathbf{x}, \mathbf{Q}^2) \otimes f_{i}^{\gamma}(\mathbf{x}_{\gamma}, \mu^2) \otimes f_{j}^{\mathrm{p}}(\mathbf{x}_{\mathrm{p}}, \mu^2) \otimes \mathrm{d}\sigma^{ij \to kl}$

- Sample x and Q^2 , setup γp sub-system with $W_{\gamma p}$
- Evolve γp as any hadronic collision (including MPIs)

Photon flux from EPA

$$f_{\gamma}^{e}(x,Q^{2}) = \frac{\alpha_{em}}{2\pi} \frac{1}{Q^{2}} \frac{(1+(1-x)^{2})^{2}}{x}$$



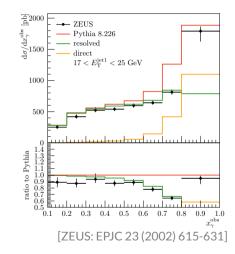
Comparison to HERA dijet photoproduction data

ZEUS dijet measurement

- $Q^2 < 1.0 \, {\rm GeV}^2$
- $134 < W_{\gamma p} < 277 \, \text{GeV}$
- $E_T^{jet1} > 14 \, GeV$, $E_T^{jet2} > 11 \, GeV$
- $-1 < \eta^{jet1,2} < 2.4$

Two contributions

- Momentum fraction of partons in photon $x_{\gamma}^{obs} = \frac{E_{T}^{jet1}e^{\eta^{jet1}} + E_{T}^{jet2}e^{\eta^{jet2}}}{2yE_{e}} \approx x_{\gamma}$
- Sensitivity to process type
- At high- $x_{\gamma}^{\rm obs}$ direct processes dominate



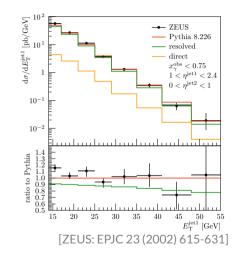
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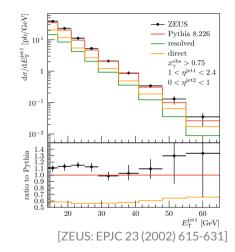
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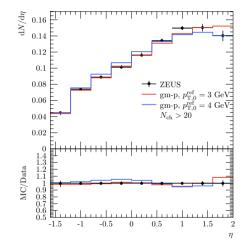
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Pseudorapidity

- Data well reproduced
- Not sensitive to MPI modelling $(p_{T,0})$



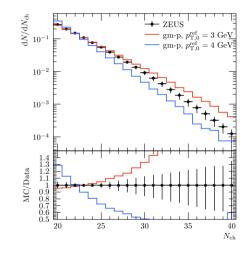
[ZEUS: JHEP 12 (2021) 102]

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Multiplicity

- Sensitivity to MPI parameters, clear support for MPIs
- Data within $p_{T,0}$ variations



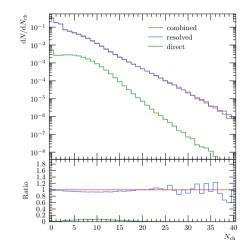
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Multiplicity

- Sensitivity to MPI parameters, clear support for MPIs
- Data within $p_{T,0}$ variations
- Direct contribution negligible in high-multiplicity events (N_{ch} > 20)



[ZEUS: JHEP 12 (2021) 102]

Alternative VMD-based approach

- Resolved contribution dominates total cross section
- \Rightarrow Set up an explicit VMD model with linear combination of vector-meson states (ρ, ω, ϕ and J/ψ)
 - Use VM PDFs from SU21

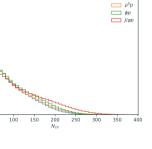
[Siöstrand, Utheim; EPJC 82 (2022) 1, 21]

Cross sections from SaS

[Schuler, Siöstrand: PRD 49 (1994) 2257-2267]

- Sample collision energy from flux
- \Rightarrow Vector meson-proton scatterings





Charged multiplicity (non-diffractive events)

2.00

1.75

1.50

1.25

△ 1.00 0.75

0.50

0.25 0.00

50

[with Marius Utheim]

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Alternative VMD-based approach

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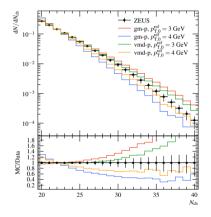
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• Cross sections from SaS

[Schuler, Sjöstrand; PRD 49 (1994) 2257-2267]

- Sample collision energy from flux
- \Rightarrow Vector meson-proton scatterings
 - In line with the full photoproduction



[ZEUS: JHEP 12 (2021) 102]

[with Marius Utheim]

Ultraperipheral collisions (UPCs)

Ultraperipheral heavy-ion collisions

- Large impact parameter (b ≥ 2R_A) ⇒ No strong interactions
- At LHC relevant for p+p, p+Pb, Pb+Pb
- Large flux due to large EM charge of nuclei
- $\Rightarrow \gamma\gamma$ and γ A collisions

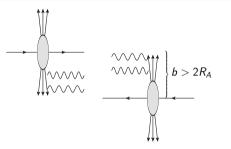
Photon flux from equivalent photon approximation

- Define flux in impact-parameter space \Rightarrow Reject hadronic interactions with b_{\min}
- Integrating the point-like approximation we get

$$f_{\gamma}^{A}(x) = \frac{2\alpha_{\rm EM}Z^{2}}{x \pi} \left[\xi \, K_{1}(\xi) K_{0}(\xi) - \frac{\xi^{2}}{2} \left(K_{1}^{2}(\xi) - K_{0}^{2}(\xi) \right) \right]$$

where $\xi = b_{\min} x m$ where $b_{\min} \approx 2R_A$ and m per nucleon mass

• Nuclear form factor heavily suppresses Q^2 of the photon \Rightarrow Photoproduction! 13



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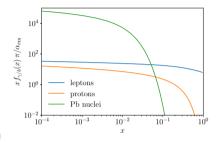


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Experimental heavy-ion UPC classification

- Event selection typically relies on Zero-degree calorimeters (X > 0)
- XnXn: At least one neutron on both sides
 - \Rightarrow A+A (hadronic interaction)
- XnOn: At least one neutron only on one side

 $\Rightarrow \gamma$ +A

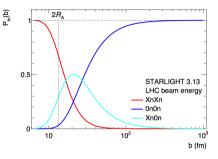
OnOn: No neutrons on either side

 $\Rightarrow \gamma + \gamma$

Possible caveats

- Additional EM interactions may break up the nuclei in "near-encounter" events
- Also diffractive processes will keep nuclei intact
 - \Rightarrow XnOn condition will remove diffractive contribution to $\gamma \text{+} \text{A}$

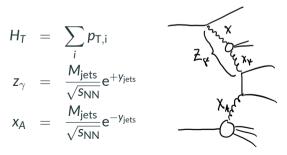




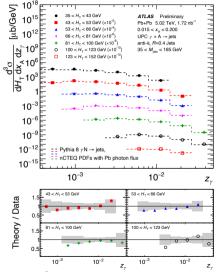
Ann.Rev.Nucl.Part.Sci. 70 (2020) 323-354

Dijets in ultra-peripheral heavy-ion collisions in Xn0n

- Good agreement out of the box when accounting both direct and resolved
- EM nuclear break-up significant
- Pythia setup with nucleon target only
 ⇒ Is such a setup enough for γ+A?



See also [Eskola, Guzey, IH, Paakkinen, Paukkunen; arXiv:2404.09731]

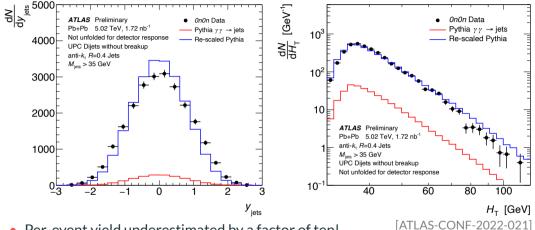


ATLAS-

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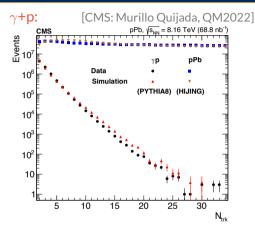
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Dijets in ultra-peripheral heavy-ion collisions in 0n0n

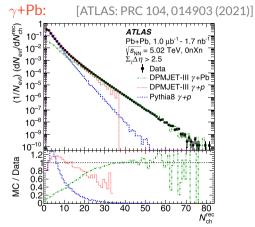


- Per-event yield underestimated by a factor of ten!
- Shape in a reasonable agreement
- $\gamma\gamma \rightarrow \mu^+\mu^-$ ok so likely a QCD effect \Rightarrow Contribution from diffractive events?

Multiplicity distributions in UPCs

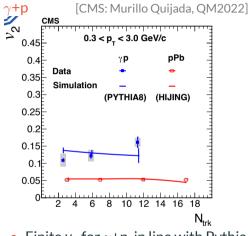


 Multiplicity distribution well reproduced in γ+p interactions



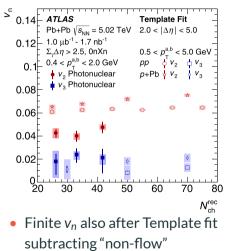
High multiplicities missed with γ+p
 ⇒ Multi-nucleon interactions

Collectivity in UPCs at the LHC



Finite v₂ for γ+p, in line with Pythia ⇒ Jet-like correlations?

γ+Pb [ATLAS: PRC 104, 014903 (2021)]



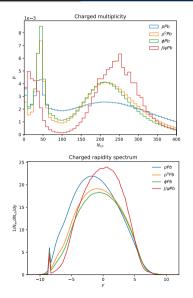
Modelling $\gamma\text{+}A$ with Pythia

[with Marius Utheim]

Angantyr model for heavy ions in Pythia

[Bierlich, Gustafson, Lönnblad, Shah; JHEP 10 (2018) 134]

- Monte Carlo Glauber to sample nucleon configurations
- Cross section fluctuations, fitted to partial nucleon-nucleon cross sections
- Secondary (wounded) collisions as diffractive excitations
- Can now handle generic hadron-ion and varying energy [I.H., Utheim; in progress]
- \Rightarrow VMD-nucleus scatterings



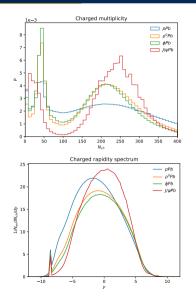
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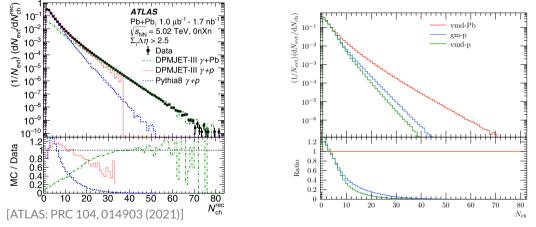
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- \Rightarrow VMD-nucleus scatterings
- \Rightarrow Hadronic cascades from cosmic rays

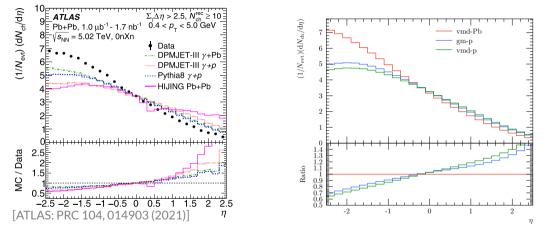


Comparison with data for γ +A



- ATLAS data not corrected for efficiency, estimated with $N_{
 m ch}^{
 m rec} pprox 0.8 \cdot N_{
 m ch}$
- Relative increase in multiplicity well in line with the VMD-Pb setup

Comparison with data for γ +A



- Multiplicity cut adjusted according to the limited efficiency
- Good description of the measured rapidity distribution with the VMD-Pb setup

Two-particle correlations in ATLAS analysis

- ATLAS apply template-fitting method to extract v_n from two-particle correlations
 - Perform a Fourier fit to obtain c_n's for low-multiplicity events (non-flow?)

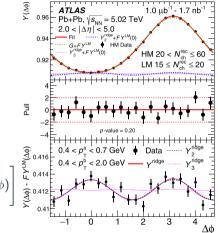
$$Y^{\text{LM}}(\Delta \phi) = c_0 + 2 \cdot \sum_{n=1}^{4} c_n \cos(n\Delta \phi)$$

• Fit high multiplicity $v_{n,n}$'s on top

$$Y^{\mathsf{HM}}(\Delta\phi) = F \cdot Y^{\mathsf{LM}}(\Delta\phi) + G\left[1 + 2 \cdot \sum_{n=2}^{4} v_{n,n} \cos(n\Delta\phi)\right]$$

Free parameters c_n , $v_{n,n}$, F, G

• Can now repeat the fit with Pythia results

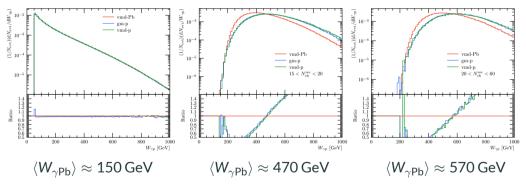


[ATLAS: PRC 104, 014903 (2021)]

Invariant mass distribution vs. multiplicity

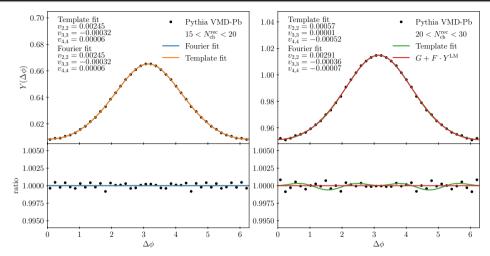
Event multiplicity increases with the collision energy

 \Rightarrow Multiplicity binning will bias W distribution



- \Rightarrow Low- and high-multiplicity event samples have different $\langle W \rangle$
- \Rightarrow Is the non-flow subtraction still meaningful?

Template fit to Pythia simulations

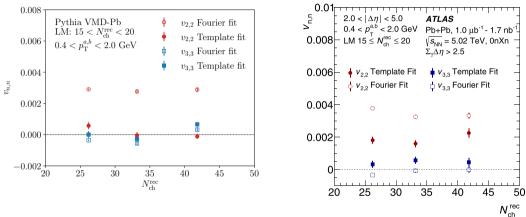


• No significant additional modulation in high-multiplicity events

Template fit = comined LM & HM fit

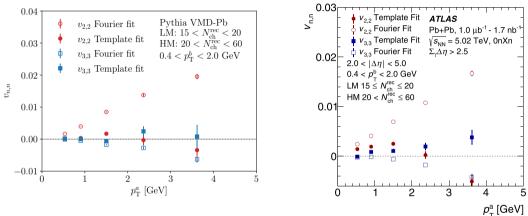
Fourier fit = direct F-fit separately to LM & HM samples ²⁴

Comparison to ATLAS $v_{n,n}$ data



- Simulated results in line with the direct Fourier fit for v_{2,2}
- Consistent with zero after template fitting (non-flow subraction)

Comparison to ATLAS $v_{n,n}$ data



- Simulated results in line with the direct Fourier fit for v_{2,2}
- Consistent with zero after template fitting (non-flow subraction)
- String interactions in high-multiplicity hadronization, hadronic rescattering?

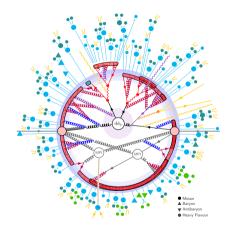
Summary & Outlook

Summary

- Ultraperipheral collisions connect LHC physics to e+p/A in HERA and EIC
- Recent HERA data will provide further constraints on Pythia implementation
- First steps towards γ +A in 8.311
 - \Rightarrow In line with multiplicity distributions
 - \Rightarrow As such not consistent with finite v₂

Outlook

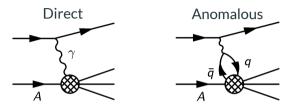
- Study different string-interaction effects for high-multiplicity events
- Study role of diffraction in UPCs

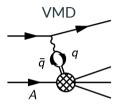


[figure by P. Skands]

Backup slides

Vector meson dominance (VMD)





Linear combination of three components

$$|\gamma
angle = c_{\mathsf{dir}}|\gamma_{\mathsf{dir}}
angle + \sum_{q} c_{q}|q\overline{q}
angle + \sum_{\mathsf{V}} c_{\mathsf{V}}|\mathsf{V}
angle$$

where the last term includes a linear combination of vector meson states up to J/Ψ

$$c_{\rm V} = \frac{4\pi\alpha_{\rm EM}}{f_{\rm V}^2}$$

V	$f_V^2/(4\pi)$
$ ho^{0}$	2.20
ω	23.6
ϕ	18.4
J/Ψ	11.5

Photon fluxes from Equivalent Photon Approximation (EPA)

• In case of a point-like lepton we have (neglecting electron mass)

$$f_{\gamma}^{I}(x,Q^{2}) = rac{lpha_{
m em}}{2\pi} rac{1}{Q^{2}} rac{(1+(1-x)^{2})}{x}$$

For protons need to include form factors, using dipole form factor

$$f_{\gamma}^{p}(x,Q^{2}) = \frac{\alpha_{\text{em}}}{2\pi} \frac{x}{Q^{2}} \frac{1}{(1+Q^{2}/Q_{0}^{2})^{4}} \left[\frac{2(1+\mu_{p}\tau)}{1+\tau} \left(\frac{1-x}{x^{2}} - \frac{M_{p}^{2}}{Q^{2}} \right) + \mu_{p}^{2} \right]$$

where $\tau = Q^2/4M_p^2$, $\mu_p = 2.79$, $Q_0^2 = 0.71 \,\text{GeV}^2$

• Drees-Zeppenfeld approximation ($M_p = 0, \mu_p = 1$)

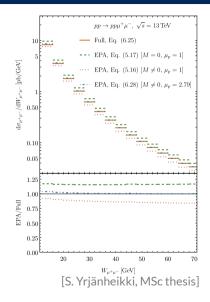
$$f_{\gamma}^{p}(x,Q^{2}) = rac{lpha_{em}}{2\pi} rac{1}{Q^{2}} rac{1}{(1+Q^{2}/Q_{0}^{2})^{4}} rac{(1+(1-x)^{2})}{x}$$

- \Rightarrow Large Q² suppressed wrt. leptons \Rightarrow photoproduction
- In ME generators (such as MG5) integrated over Q² and assumed collinear

Equivalent photon approximation

Compare to full calculation

- Example process $pp \rightarrow \gamma \gamma \rightarrow \mu^+ \mu^-$
- Different approximations (e.g.) by Drees and Zeppenfeld \sim 20% difference to full calculation
- Keeping finite mass and correct magnetic moment provides ~ few percent accuracy
- Not checked for other observables, such as acoplanarity



Define your own photon flux for Pythia 8

• Derive a new object from PDF class

class Proton2gammaEPA : public PDF {

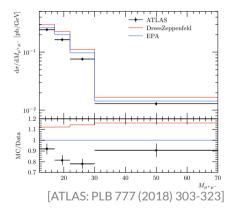
public:

```
// Constructor.
ProtonZgammaEPA(int idBeamIn) : PDF(idBeamIn) {}
// Update the photon flux.
void xfUpdate(int , double x, double Q2) {
    double mzproton = pow2(20.938);
    double mzproton = pow2(2.79);
    double Q20 = 0.71;
    double Q20 = 0.71;
    double Q20 = 0.71;
    double Q20 = 0.71;
    double Coupling = 0.5 * 0.007297353080 / N_PI * FQ4;
    double tou = Q2 / (4. * m2proton);
    xgamma = coupling * ( pow2(x) / Q2 ) * ( 2. * (1. + mup2*tau ) / (1. + tau)
        * ( (1 - x)/pow2(x) - m2proton / Q2 ) + mup2);
}.
```

• Pass as a pointer to Pythia

pythia.readString("PDF:becm&2gamma = on"); pythia.readString("PDF:becm&2gamma5et = 0"); pythia.readString("PDF:proton2gamma5et = 0"); PDFPtr photonFluxA = make_shared<Proton2gammaEPA>(2212); PDFPtr photonFluxB = make_shared<Proton2gammaEPA>(2212); pythia.setPhotonFluxBtr(photonFluxA, photonFluxB);

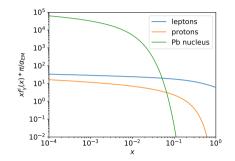
Example in p-p: $\gamma\gamma \rightarrow \mu^+\mu^-$



• No finite-size effects accounted

• Enable γ +p in e+p

pythia.readString("Beams:idA = -11"); pythia.readString("Beams:idB = 2212"); pythia.readString("PDF:beamA2gamma = on");

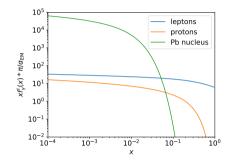


• Enable γ +p in e+p

pythia.readString("Beams:idA = -11");
pythia.readString("Beams:idB = 2212");
pythia.readString("PDF:beamA2gamma = on");

• Enable γ +p in p+p

pythia.readString("Beams:idA = 2212");
pythia.readString("Beams:idB = 2212");
pythia.readString("PDF:beamA2gamma = on");



• Enable γ +p in e+p

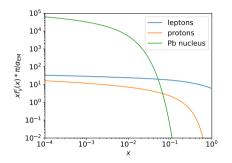
pythia.readString("Beams:idA = -11");
pythia.readString("Beams:idB = 2212");
pythia.readString("PDF:beamA2gamma = on");

• Enable γ +p in p+p

pythia.readString("Beams:idA = 2212");
pythia.readString("Beams:idB = 2212");
pythia.readString("PDF:beamA2gamma = on");

• Enable γ +p in Pb+p

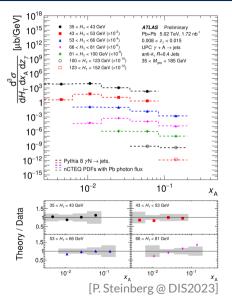
pythia.readString("Beams:idA = 2212"); pythia.readString("PDF:beamA2gamma = on"); pythia.readString("PDF:beamA2gammaSet = 0"); pythia.readString("PDF:beam2gammaApprox = 2"); pythia.readString("Photon:sample02 = off"); PDFPtr photonFlux = make_shared<Nucleus2gamma>(2212); pythia.setPhotonFluxPtr(photonFlux, 0);



For more examples see main68.cc,main69.cc, main70.cc,main78.cc in examples directory



[from main70.cc]

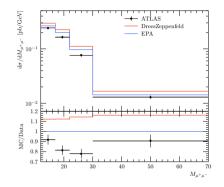


An example process: $\gamma\gamma \rightarrow \mu^+\mu^-$

- Can take place in EE, SD and DD (also DY processes with resolved photons?)
- Implemented natively in Pythia, can also generate with an ME generator (MG5, SC)

EE contribution

- Clean process to study fluxes
- However, fluxes only does not account for finite-size effects



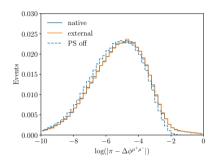
[ATLAS: PLB 777 (2018) 303-323]

An example process: $\gamma\gamma \rightarrow \mu^+\mu^-$

- Can take place in EE, SD and DD (also DY processes with resolved photons?)
- Implemented natively in Pythia, can also generate with an ME generator (MG5, SC)

EE contribution

- Clean process to study fluxes
- However, fluxes only does not account for finite-size effects
- Not quite back-to-back due to
 - *p*_T generated by non-collinear photons
 - QED radiation in the final state
- Acoplanarity $|\pi \Delta \phi|$ quantify the effect



- Needed to tune Pythia primordial k_T parameters for external events
- Can use (user-defined) flux for Q² sampling

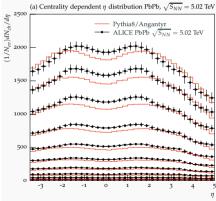
Heavy-ion collisions

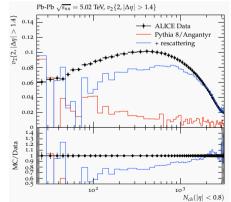
• Angantyr in Pythia provides a full heavy-ion collisions framework

[Bierlich, Gustafson, Lönnblad & Shah: 1806.10820]

• Hadronic rescattering can be included as well, enhances collective effects

[CB, Ferreres-Solé, Sjöstrand & Utheim: 1808.04619, 2005.05658, 2103.09665]

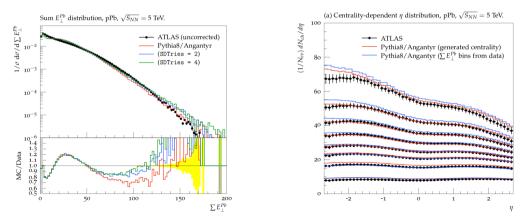




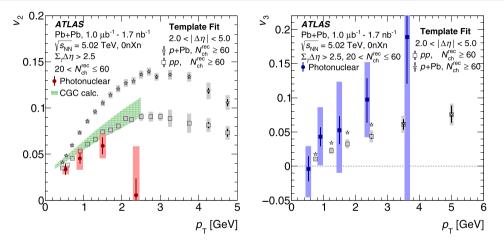
p+A collisions

[Bierlich, Gustafson, Lönnblad & Shah: 1806.10820]

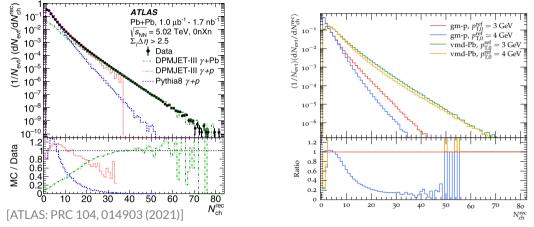
- Angantyr can be applied also to asymmetric p+A collisions
- The centrality measure well reproduced
- Similarly centraility-dependent multiplicities



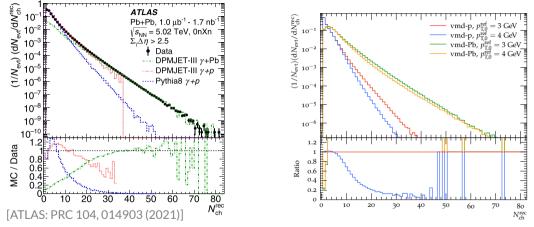
ATLAS data for v_n in γ +Pb



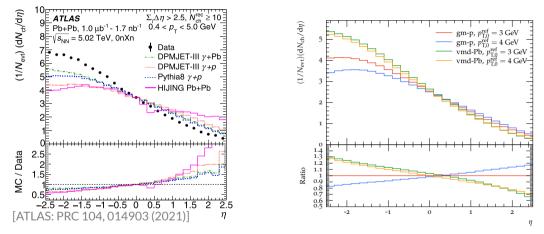
- Non-zero flow coefficients also for γ+Pb
- Expected baseline from MC simulations?



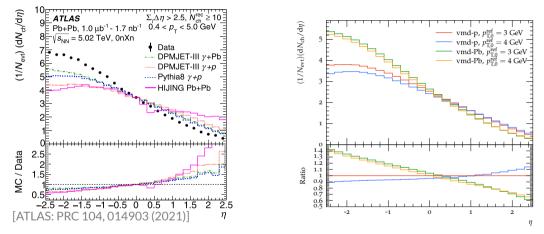
- Pythia8 γ +p in ATLAS result should correspond to gm-p on right
- Relative increase in multiplicity well in line with the VMD setup



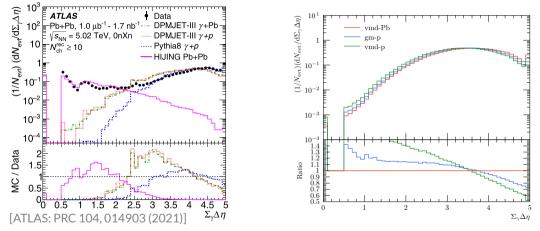
- Pythia8 γ +p in ATLAS result should correspond to gm-p on right
- Relative increase in multiplicity well in line with the VMD setup



- Pythia8 γ+p in ATLAS result should correspond to gm-p on right
- Relative shift in rapidity distribution in line with the VMD setup using Angantyr

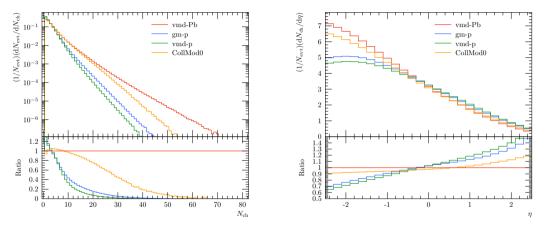


- Pythia8 γ +p in ATLAS result should correspond to gm-p on right
- Relative shift in rapidity distribution in line with the VMD setup using Angantyr



- $\Sigma_{\gamma} \Delta \eta$: Sum of rapidity gaps for which $\Delta \eta > 0.5$
- Similar for γ -p and γ -Pb

Role of cross section fluctuations



• High-multiplicity tail less pronounced with Angantyr:CollisionModel = 0 with fixed nucleon radius, ATLAS data seem to favour fluctuations